TANGENT DIRAC STRUCTURES OF HIGHER ORDER

P. M. KOUOTCHOP WAMBA, A. NTYAM, AND J. WOUAFO KAMGA

ABSTRACT. Let L be an almost Dirac structure on a manifold M. In [2] Theodore James Courant defines the tangent lifting of L on TM and proves that:

If L is integrable then the tangent lift is also integrable.

In this paper, we generalize this lifting to tangent bundle of higher order.

INTRODUCTION

Let M be a differential manifold (dim M = m > 0). Consider the mapping ϕ_M defined by:

$$\phi_M : \quad TM \oplus T^*M \times_M TM \oplus T^*M \quad \to \quad \mathbb{R}$$
$$((X_1, \alpha_1), (X_2, \alpha_2)) \qquad \mapsto \quad \frac{1}{2} (\langle X_1, \alpha_2 \rangle_M + \langle X_2, \alpha_1 \rangle_M)$$

where $\langle \cdot \rangle_M$ is the canonical pairing defined by:

$$\begin{array}{rccc} TM \times_M T^*M & \to & \mathbb{R} \\ (X, \alpha) & \mapsto & \langle X, \alpha \rangle_M \end{array}$$

An almost Dirac structure on M, is a sub vector bundle L of the vector bundle $TM \oplus T^*M$, which is isotropic with respect to the natural indefinite symmetric scalar product ϕ_M (i.e $\forall (X_1, \alpha_1), (X_2, \alpha_2) \in \Gamma(L), \phi_M((X_1, \alpha_1), (X_2, \alpha_2)) = 0$), and such that the rank of L is equal to the dimension of M.

We define on the set $\Gamma(TM \oplus T^*M)$ of sections of $TM \oplus T^*M$ a bracket by:

$$\forall (X_1, \alpha_1), (X_2, \alpha_2) \in \Gamma(TM \oplus T^*M) [(X_1, \alpha_1), (X_2, \alpha_2)]_C = ([X_1, X_2], \mathcal{L}_{X_1}\alpha_2 - i_{X_2}d\alpha_1).$$

This bracket is called Courant bracket. A Dirac structure (or generalized Dirac structure) is an almost Dirac structure such that:

$$\forall (X_1, \alpha_1), (X_2, \alpha_2) \in \Gamma(L), \qquad [(X_1, \alpha_1), (X_2, \alpha_2)] \in \Gamma(L).$$

This condition is called "integrability condition".

²⁰¹⁰ Mathematics Subject Classification: primary 53C15; secondary 53C75, 53D05.

Key words and phrases: Dirac structure, almost Dirac structure, tangent functor of higher order, natural transformations.

Received August 24, 2009, revised August 2010. Editor I. Kolář.

For $(X_3, \alpha_3) \in \Gamma(TM \oplus T^*M)$, in [2] is defined the 3-tensor $T_{TM \oplus T^*M}$ on the vector bundle $TM \oplus T^*M$ by:

 $T_{TM\oplus T^*M}((X_1,\alpha_1),(X_2,\alpha_2),(X_3,\alpha_3)) = \phi_M\big([(X_1,\alpha_1),(X_2,\alpha_2)],(X_3,\alpha_3)\big).$

We put $T_L = T_{TM \oplus T^*M}|_{\Gamma(L) \times \Gamma(L) \times \Gamma(L)}$. The integrability condition of L is determined by the vanishing of the 3-tensor T_L on the vector bundle L.

For all integer $r, k \ge 1$, we have the jet functor T_k^r of k-dimensional velocity of order r and, when k = 1, this functor is denoted by T^r and is called tangent bundle of order r. When r = 1, T^1 is a natural equivalence of tangent functor T.

The main results of this paper are theorems 2 and 3: giving an almost Dirac structure L on M, we construct an almost Dirac structure L^r on T^rM and we prove that: L is integrable if and only if L^r is integrable.

All manifolds and maps are assumed to be infinitely differentiable. r will be a natural integer $(r \ge 1)$.

1. Other characterization of generalized Dirac structure

Let V be a real vector space of dimension m. We consider the map

$$\begin{array}{rcl} \phi_V \colon & V \oplus V^* \times V \oplus V^* & \to & \mathbb{R} \\ & & \left((u, u^*), (v, v^*) \right) & \mapsto & \frac{1}{2} \big(\langle u, v^* \rangle + \langle v, u^* \rangle \big) \end{array}$$

where $\langle \cdot \rangle$ is the dual bracket $V \times V^* \to \mathbb{R}$.

Definition 1. A constant Dirac structure on V is a sub vector space L of dimension m of $V \oplus V^*$ such that:

$$\forall (u, u^*), (v, v^*) \in L, \qquad \phi_V((u, u^*), (v, v^*)) = 0.$$

Theorem 1. A constant Dirac structure L on V is determined by a pair of linear maps $a: \mathbb{R}^m \to V$ and $b: \mathbb{R}^m \to V^*$ such that:

$$(1) a^* \circ b + b^* \circ a = 0$$

$$(2) \qquad \ker a \cap \ker b = \{0\}$$

Proof. Condition (1) is the isotropy of constant Dirac structure, and condition (2) is the maximality of the isotropy. \Box

Remark 1.

- (1) We say that the constant Dirac structure L is determined by the linear maps a and b.
- (2) An almost Dirac structure on a differential manifold M is a sub vector bundle of $TM \oplus T^*M$ such that: $\forall x \in M$, the fiber L_x of L over x is a constant Dirac structure on T_xM .
- (3) An almost Dirac structure at a point $x \in M$ is determined by a pair of maps $a_x : \mathbb{R}^m \to T_x M, b_x : \mathbb{R}^m \to T_x^* M$ such that:

$$\begin{cases} a_x^* \circ b_x + b_x^* \circ a_x = 0\\ \ker a_x \cap \ker b_x = \{0\} \end{cases}$$

Corollary. An almost Dirac structure is determined in a neighbourhood U of a local trivialization $L|_U \approx U \times \mathbb{R}^m$ by a pair of vector bundle morphisms $a: U \times \mathbb{R}^m \to T_U M$, $b: U \times \mathbb{R}^m \to T_U^* M$ over U such that:

$$\forall x \in U, \qquad \begin{cases} a_x^* \circ b_x + b_x^* \circ a_x = 0\\ \ker a_x \cap \ker b_x = \{0\} \end{cases}$$

We denote by p_1 and p_2 the natural projections of $TM \oplus T^*M$ onto TM and T^*M respectively. Note that $a: L \to TM$ and $b: L \to T^*M$ are really globally defined and are nothing more than the projections p_1 and p_2 .

Example 1. Let M be an m-dimensional manifold.

(1) Let ω be a differential form on M of degree 2.

$$\Gamma = \{ (X, i_X \omega), \quad X \in \mathfrak{X}(M) \}.$$

 Γ is the set of differential sections of an almost Dirac structure on M. It is a Dirac structure if and only if ω is pre-symplectic form.

(2) Let Π be a bivector field on M.

$$\Gamma' = \{ (i_{\Pi}\alpha, \alpha), \quad \alpha \in \Omega^1(M) \}.$$

 Γ' is the set of differential sections of an almost Dirac structure on M. It is a Dirac structure if and only if Π is a Poisson bivector.

We denote by (x^i, \dot{x}^i) and (x^i, p_i) a local coordinates system of TM and T^*M respectively. Let L be an almost Dirac structure on M defined locally by:

$$a: U \times \mathbb{R}^m \to TM$$
 and $b: U \times \mathbb{R}^m \to T^*M$.

We have:

$$\begin{cases} a(x^i, e_j) = a_j^k \frac{\partial}{\partial x^k} \\ b(x^i, e_j) = b_{jk} dx^k \end{cases}$$

where (e_i) denote the canonical basis of \mathbb{R}^m . Locally the 3-tensor field T_L is:

$$T_L = \sum_{\text{cyclic},i,j,k} \left(a_i^p \frac{\partial b_{js}}{\partial x^p} a_k^s + a_i^p \frac{\partial a_j^s}{\partial x^p} b_{ks} \right).$$

2. TANGENT DIRAC STRUCTURE OF HIGHER ORDER

 $\kappa_M^r: T^rTM \to TT^rM$ and $\alpha_M^r: T^*T^rM \to T^rT^*M$ denote the natural transformations defined in [1] and [7]. We have:

$$\langle \kappa_M^r(u), v^* \rangle_{T^r M} = \langle u, \alpha_M^r(v^*) \rangle_{T^r M}', \qquad (u, v^*) \in T^r T M \times_{T^r M} T^* T^r M$$

where $\langle \cdot \rangle'_{T^r M} = \tau_r \circ T^r \langle \cdot \rangle$ and $\tau_r(j_0^r \varphi) = \frac{-\tau}{dt^r}(t)|_{t=0}$. We denote by ε_M^r the inverse map of α_M^r .

Consider the maps $a: U \times \mathbb{R}^m \to TM$ and $b: U \times \mathbb{R}^m \to T^*M$. We take their tangents of order r, to get:

$$T^r a \colon T^r U \times \mathbb{R}^{m(r+1)} \to T^r T M$$
 and $T^r b \colon T^r U \times \mathbb{R}^{m(r+1)} \to T^r T^* M$.

We apply natural transformations κ_M^r and ε_M^r respectively, to get the vector bundle maps over id_{T^rU} defined by:

 $a^r \colon T^r U \times \mathbb{R}^{m(r+1)} \to T T^r M$ and $b^r \colon T^r U \times \mathbb{R}^{m(r+1)} \to T^* T^r M$.

Theorem 2. The pair of maps a^r and b^r determines a generalized almost Dirac structure L^r on T^rM , which we call the tangent lift of order r of the generalized almost Dirac structure on M determined by a and b.

Proof. Firstly, we prove that: $(a^r)^* \circ b^r + (b^r)^* \circ a^r = 0$. Let $j_0^r \psi, j_0^r \varphi \in T^r(U \times \mathbb{R}^m)$, where $\varphi, \psi \colon \mathbb{R} \to U \times \mathbb{R}^m$ differentials. We have:

$$\begin{split} \langle (a^r)^* \circ b^r(j_0^r \varphi), j_0^r \psi \rangle &= \langle b^r(j_0^r \varphi), a^r(j_0^r \psi) \rangle \\ &= \langle \varepsilon_M^r \circ T^r b, \kappa_M^r \circ T^r a(j_0^r \psi) \rangle \\ &= \langle T^r b(j_0^r \varphi), T^r a(j_0^r \psi) \rangle'_{T^r M} \\ &= \tau^r \circ j_0^r(\langle b \circ \varphi, a \circ \psi \rangle_M) \\ &= \tau^r \circ j_0^r(\langle a^* \circ b \circ \varphi, \psi \rangle_M) \,. \end{split}$$

By the same way, we have:

$$\langle (b^r)^* \circ a(j_0^r \varphi), j_0^r \psi \rangle = \tau^r \circ j_0^r(\langle b^* \circ a \circ \varphi, \psi \rangle_M)$$

we deduce that:

$$\left\langle ((a^r)^* \circ b^r + (b^r)^* \circ a)(j_0^r \varphi), j_0^r \psi \right\rangle = \tau^r \circ j_0^r \left(\left\langle (a^* \circ b + b^* \circ a) \circ \varphi, \psi \right\rangle_M \right) = 0.$$

Secondly we prove that: ker $a^r \cap \ker b^r = \{0\}$. We prove this case for r = 2. The proof for $r \ge 3$ is similar.

In the local coordinates system, we have:

$$a: U \times \mathbb{R}^{m} \to U \times \mathbb{R}^{m} \text{ and } b: U \times \mathbb{R}^{m} \to U \times (\mathbb{R}^{m})^{*}$$
$$(x, e) \mapsto (x, ae) \text{ and } (x, e) \mapsto (x, be)$$
$$a^{2}(x, \dot{x}, \ddot{x}, e, \dot{e}, \ddot{e}) = (x, \dot{x}, \ddot{x}, ae, \dot{a}e + a\dot{e}, \ddot{a}e + \dot{a}\dot{e} + a\ddot{a})$$
$$b^{2}(x, \dot{x}, \ddot{x}, e, \dot{e}, \ddot{e}) = (x, \dot{x}, \ddot{x}, ae, \dot{a}e + a\dot{e}, \ddot{a}e + \dot{a}\dot{e} + a\ddot{a})$$
$$b^{2}(x, \dot{x}, \ddot{x}, e, \dot{e}, \ddot{e}) = (x, \dot{x}, \ddot{x}, be + b\dot{e} + b\ddot{e}, be + b\dot{e}, be)$$
$$a^{2}(e, \dot{e}, \ddot{e}) = \begin{pmatrix} a & 0 & 0 \\ \dot{a} & a & 0 \\ \ddot{a} & \dot{a} & a \end{pmatrix} \begin{pmatrix} e \\ \dot{e} \\ \ddot{e} \end{pmatrix} \text{ and } b^{2}(e, \dot{e}, \ddot{e}) = \begin{pmatrix} \ddot{b} & \dot{b} & b \\ \dot{b} & b & 0 \\ b & 0 & 0 \end{pmatrix} \begin{pmatrix} e \\ \dot{e} \\ \ddot{e} \end{pmatrix}.$$

$$a^{2}(e, \dot{e}, e) = \begin{pmatrix} a & a & 0 \\ \ddot{a} & \dot{a} & a \end{pmatrix} \begin{pmatrix} e \\ \ddot{e} \end{pmatrix} \quad \text{and} \quad b^{2}(e, e, e) = \begin{pmatrix} b & 0 \\ b & 0 \end{pmatrix}$$
$$a^{2}(e, \dot{e}, \ddot{e}) = b^{2}(e, \dot{e}, \ddot{e}) = 0 \text{ we have:}$$

If $a^{2}(e, \dot{e}, \ddot{e}) = b^{2}(e, \dot{e}, \ddot{e}) = 0$, we have:

$$ae = 0$$
 $be = 0$ \Rightarrow $e \in \ker a \cap \ker b = \{0\}$.

and it follows that e = 0.

$$\begin{cases} b\dot{e} + \dot{b}e = 0\\ a\dot{e} + \dot{a}e = 0 \end{cases} \Rightarrow \begin{cases} b\dot{e} = 0\\ a\dot{e} = 0 \end{cases}$$

e and \dot{e} are constant, it follows that $\dot{e} = 0$.

$$\begin{cases} b\ddot{e} = 0\\ a\ddot{e} = 0 \end{cases} \quad \Rightarrow \quad \ddot{e} = 0 \,.$$

Thus ker $a^2 \cap \ker b^2 = \{0\}.$

Theorem 3. The almost Dirac structure L on M is integrable if and only if the almost Dirac structure L^r on T^rM is integrable.

Proof. Consider the local coordinates system $\{x^1, \ldots, x^m\}$ of M, we have:

$$a(x^i, e_j) = a_k^i \frac{\partial}{\partial x^k}$$
 and $b(x^i, e_j) = b_{ik} dx^k$

We have:

$$a^{r} = \begin{pmatrix} a_{j}^{i} & \dots & 0\\ \vdots & \dots & \vdots\\ a_{j}^{(r)} & & \\ a_{j}^{i} & \dots & a_{j}^{i} \end{pmatrix} \quad \text{and} \quad b^{r} = \begin{pmatrix} (r) & & & \\ b_{ij} & \dots & b_{ij}\\ \vdots & & & \vdots\\ b_{ij} & \dots & 0 \end{pmatrix}$$

We get $a^r = (A_j^i)_{1 \le i,j \le m(r+1)}$ and $b^r = (B_{ij})_{1 \le i,j \le m(r+1)}$. For q, d = 0, 1, ..., r, we have:

$$\begin{aligned} \forall (i,j) \in & \{qm+1,\ldots,m(q+1)\} \times \{dm+1,\ldots,m(d+1)\}, \\ & \left\{ \begin{aligned} A_j^i &= (a_{j-md}^{i-mq})^{(q-d)} \\ & \text{and} \\ B_{ij} &= (b_{i-mq,j-md})^{(r-q-d)} \end{aligned} \right. \end{aligned}$$

We adopt the following notation:

$$\frac{\partial}{\partial x^p} = \frac{\partial}{\partial x^{p-m\alpha}_{\alpha}} = \left(\frac{\partial}{\partial x^{p-m\alpha}}\right)^{(\alpha)} \quad \left(\alpha m + 1 \le p \le \alpha(m+1)\right).$$

The Courant tensor T_{ijk} of the almost Dirac structure is given by:

$$T_{ijk} = \sum_{\text{cyclic}, i, j, k} A_i^p \frac{\partial B_{js}}{\partial x^p} A_k^s + A_i^p \frac{\partial A_j^s}{\partial x^p} B_{ks}, \quad \text{we wish to verify that} \quad T_{ijk} = 0.$$

We take $hm + 1 \le i \le m(h+1)$, $\ell m + 1 \le j \le m(\ell+1)$ and $tm + 1 \le k \le m(t+1)$ for $h, \ell, t = 0, 1, ..., r$. We have:

$$T_{ijk} = \sum_{q=0}^{r} \sum_{d=0}^{r} \sum_{p=qm+1}^{r} \sum_{s=dm+1}^{d(m+1)} \left(A_i^p \frac{\partial B_{js}}{\partial x^p} A_k^s + A_i^p \frac{\partial A_j^s}{\partial x^p} B_{ks} \right)$$
$$= (a_{i-mh}^{p-mq})^{(q-h)} \frac{\partial (b_{j-m\ell,s-md})^{(r-\ell-d)}}{\partial x_q^{p-mq}} (a_{k-mt}^{s-md})^{(d-\ell)} + (a_{i-mh}^{p-mq})^{(q-h)} \frac{\partial (a_{j-m\ell}^{s-md})^{(d-\ell)}}{\partial x_q^{p-mq}} (b_{k-mt,s-md})^{(r-d-t)}$$

$$= (a_{i-mh}^{p-mq})^{(q-h)} \left(\frac{\partial b_{j-m\ell,s-md}}{\partial x^{p-mq}}\right)^{(r-\ell-d-q)} (a_{k-mt}^{s-md})^{(d-t)} + (a_{i-mh}^{p-mq})^{(q-h)} \left(\frac{\partial a_{j-m\ell}^{s-md}}{\partial x^{p-mq}}\right)^{(d-\ell-q)} (b_{k-mt,s-md})^{(r-d-t)} = \left(a_{i-mh}^{p-md} \frac{\partial b_{j-m\ell,s-md}}{\partial x^{p-mq}} a_{k-mt}^{s-md}\right)^{(r-\ell-h-t)} + \left(a_{i-mh}^{p-mq} \frac{\partial a_{j-m\ell}^{s-md}}{\partial x^{p-mq}} b_{k-mt,s-md}\right)^{(r-\ell-h-t)} = (a_{i-mh}^{p-mq} \frac{\partial b_{j-m\ell,s-md}}{\partial x^{p-mq}} a_{k-mt}^{s-md} + a_{i-mh}^{p-mq} \frac{\partial a_{j-m\ell}^{s-md}}{\partial x^{p-mq}} b_{k-mt,s-md})^{(r-\ell-h-t)}$$

the calculation above shows that $T_L = 0$ if and only if $T_{L^r} = 0$.

Remark 2. This construction generalizes the tangent lifts of higher order of Poisson and pre-symplectic structure to tangent bundle of higher order.

 \square

References

- Cantrijn, F., Crampin, M., Sarlet, W., Saunders, D., The canonical isomorphism between T^kT^{*} and T^{*}T^k, C. R. Acad. Sci., Paris, Sér. II **309** (1989), 1509–1514.
- [2] Courant, T., Tangent Dirac Structures, J. Phys. A: Math. Gen. 23 (22) (1990), 5153–5168.
- [3] Courant, T., Tangent Lie Algebroids, J. Phys. A: Math. Gen. 27 (13) (1994), 4527–4536.
- [4] Gancarzewicz, J., Mikulski, W., Pogoda, Z., Lifts of some tensor fields and connections to product preserving functors, Nagoya Math. J. 135 (1994), 1–41.
- [5] Grabowski, J., Urbanski, P., Tangent lifts of Poisson and related structures, J. Phys. A: Math. Gen. 28 (23) (1995), 6743–6777.
- [6] Kolář, I., Michor, P., Slovák, J., Natural Operations in Differential Geometry, Springer-Verlag, 1993.
- [7] Morimoto, A., Lifting of some type of tensors fields and connections to tangent bundles of p^r-velocities, Nagoya Math. J. 40 (1970), 13–31.
- [8] Ntyam, A., Wouafo Kamga, J., New versions of curvatures and torsion formulas of complete lifting of a linear connection to Weil bundles, Ann. Pol. Math. 82 (3) (2003), 233–240.

Department of Mathematics, The University of Yaoundé 1, P.O BOX, 812, Yaoundé, Cameroon *E-mail*: wambapm@yahoo.fr

DEPARTMENT OF MATHEMATICS, ENS YAOUNDÉ, P.O BOX, 47 YAOUNDÉ, CAMEROON *E-mail*: antyam@uy1-uninet.cm

Department of Mathematics, The University of Yaoundé 1, P.O BOX, 812, Yaoundé Cameroon *E-mail*: wouafoka@yahoo.fr