# Higher order Osserman pseudo-Riemannian manifolds of neutral signature $(2,2)$ 

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#### Abstract

In this paper we construct a family of pseudo-Riemannian metrics of neutral signature $(2,2)$ which leads to $k$-Osserman manifolds for all $k$ admissible. For these manifolds the generalized Jacobi operator is 2-nilpotent. Conditions for locally symmetry on the considered manifolds are established.


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Key words: generalized Jacobi operator, locally symmetric.
Let $(M, g)$ be a pseudo-Riemannian manifold of signature $(p, q)$ and dimension $n=p+q$. Let $R(\cdot, \cdot)$ be the Riemannian curvature operator. The Jacobi operator $J(X): Y \rightarrow R(Y, X) X$ is a self-adjoint operator and it plays an important role in the study of geodesic variations.

Let $S^{ \pm}(M)$ be the pseudo-sphere bundles of unit spacelike $(+)$ and timelike $(-)$ vectors for the manifold $(M, g)$. Then $(M, g)$ is said to be spacelike Osserman (respectively timelike Osserman) if the eigenvalues of $J(\cdot)$ are constant on $S^{+}(M, g)$ (respectively on $S^{-}(M, g)$ ). The notions spacelike Osserman and timelike Osserman are equivalent and if $(M, g)$ is either of them, then $(M, g)$ is said to be Osserman.

In this paper we study the higher order Jacobi operator, which was first defined by Stanilov and Videv ([9]) in the Riemannian setting. This definition was extended to semi-Riemannian geometry in [6]. Let $\pi$ be a nondegenerate $k$-plane in $T_{p} M$, with orthonormal basis $\left\{e_{1}, \ldots, e_{k}\right\}$, where $(M, g)$ is a pseudo-Riemannian manifold of signature $(p, q)$. The generalized Jacobi operator is defined by

$$
J_{R}(\pi)=\sum_{i=1}^{k} g\left(e_{i}, e_{i}\right) R\left(\cdot, e_{i}\right) e_{i}
$$

We say that a pair of integers $(r, s)$ is an admissible pair for $T_{p} M$ if $0 \leq r \leq p$, $0 \leq s \leq q$ and $1 \leq s+r \leq p+q-1$. This means that the Grassmannian $G r_{(r, s)}\left(T_{p} M\right)$ of all non-degenerate planes in $T_{p} M$ of signature $(r, s)$ is non-empty and does not consist of a single point.

Let $(r, s)$ be an admissible pair. We say that $(M, g)$ is Ossermann of type $(r, s)$ in $p \in M$ if the eigenvalues of the operator $J_{R}(\pi)$ do not depend on the choice of plane $\pi \in G r_{(r, s)}\left(T_{p} M\right)$.
P. Gilkey shows that if $(M, g)$ is Osserman of type $(r, s)$ then it is Osserman of type $(\tilde{r}, \tilde{s})$ for all admissible pairs $(\tilde{r}, \tilde{s})$ satisfying $r+s=\tilde{r}+\tilde{s}([3],[4])$. Thus, only the dimension $k=r+s$ of planes $\pi$ is relevant and we simply talk about $k$-Osserman. A semi-Riemannian manifold $(M, g)$ is said to be a $k$-Osserman manifold if for all points $p \in M,(M, g)$ is $k$-Osserman in $p$ with the eigenvalue structure of $J_{R_{p}}(\cdot)$ independent of the chosen point $p$.

Let $M=\mathbf{R}^{4}$ with coordinates $(x, y)=\left(x^{1}, x^{2}, y^{1}, y^{2}\right)$. Then $\mathcal{X}=\operatorname{Span}\left\{\partial_{1}^{x}, \partial_{2}^{x}\right\}$ and $\mathcal{Y}=\operatorname{Span}\left\{\partial_{1}^{y}, \partial_{2}^{y}\right\}$ define two distributions of $T M$. The splitting $T M=\mathcal{X} \bigoplus \mathcal{Y}$ is just the usual splitting $T \mathbf{R}^{4}=T \mathbf{R}^{2} \bigoplus T \mathbf{R}^{2}$. We define a semi-Riemannian metric of neutral signature $(2,2)$ by setting

$$
\begin{align*}
g_{\left(f_{1}, f_{2}, h\right)} & =y^{1} f_{1}\left(x^{1}\right) d x^{1} \otimes d x^{1}+y^{2} f_{2}\left(x^{2}\right) d x^{2} \otimes d x^{2}+ \\
& +h\left(x^{1}, x^{2}\right)\left[d x^{1} \otimes d x^{2}+d x^{2} \otimes d x^{1}\right]+  \tag{0.1}\\
& +a\left[d x^{1} \otimes d y^{1}+d y^{1} \otimes d x^{1}+d x^{2} \otimes d y^{2}+d y^{2} \otimes d x^{2}\right]
\end{align*}
$$

where $a \in \mathbf{R}^{*}$ and $f_{1}, f_{2}, h$ are smooth real valued functions. The coefficients of $g_{\left(f_{1}, f_{2}, h\right)}$ depend on $x$ and $y$. Furthermore, the distribution $\mathcal{Y}$ is totally isotropic with respect to $g_{\left(f_{1}, f_{2}, h\right)}$.

Lemma 1 The only nonvanishing covariant derivatives are given by

$$
\begin{align*}
\nabla \partial_{1}^{x} \partial_{1}^{x} & =-\frac{1}{2 a} f_{1}\left(x^{1}\right) \partial_{1}^{x}+\left[\frac{1}{2 a} y^{1} f_{1}^{\prime}\left(x^{1}\right)+\frac{y^{1}}{2 a^{2}} f_{1}^{2}\left(x^{1}\right)\right] \partial_{1}^{y}+ \\
& +\left[\frac{1}{a} \frac{\partial h}{\partial x^{1}}\left(x^{1}, x^{2}\right)+\frac{1}{2 a^{2}} f_{1}\left(x^{1}\right) h\left(x^{1}, x^{2}\right)\right] \partial_{2}^{y}, \\
\nabla \partial_{2}^{x} \partial_{2}^{x} & =-\frac{1}{2 a} f_{2}\left(x^{2}\right) \partial_{2}^{x}+\left[\frac{1}{2 a^{2}} f_{2}\left(x^{2}\right) h\left(x^{1}, x^{2}\right)+\frac{1}{a} \frac{\partial h}{\partial x^{2}}\left(x^{1}, x^{2}\right)\right] \partial_{1}^{y}+  \tag{0.2}\\
& +\left[\frac{1}{2 a} y^{2} f_{2}^{\prime}\left(x^{2}\right)+\frac{y^{2}}{2 a^{2}} f_{2}^{2}\left(x^{2}\right)\right] \partial_{2}^{y}, \\
\nabla \partial_{1}^{x} \partial_{1}^{y} & =\frac{1}{2 a} f_{1} \partial_{1}^{y}, \\
\nabla \partial_{2}^{x} \partial_{2}^{y} & =\frac{1}{2 a} f_{2} \partial_{2}^{y} .
\end{align*}
$$

From (0.1) we have the following:
Proposition 1 The only nonvanishing components of the curvature tensor of $\left(\mathbf{R}^{4}, g_{\left(f_{1}, f_{2}, h\right)}\right)$ are given by

$$
\begin{align*}
& R\left(\partial_{1}^{x}, \partial_{2}^{x}\right) \partial_{1}^{x}=-\frac{1}{a}\left[\frac{\partial^{2} h}{\partial x^{1} \partial x^{2}}+\frac{1}{2 a} f_{2} \frac{\partial h}{\partial x^{1}}+\frac{1}{2 a} f_{1} \frac{\partial h}{\partial x^{2}}+\frac{1}{4 a^{2}} f_{1} f_{2} h\right] \partial_{2}^{y}  \tag{0.3}\\
& R\left(\partial_{1}^{x}, \partial_{2}^{x}\right) \partial_{2}^{x}=\frac{1}{a}\left[\frac{\partial^{2} h}{\partial x^{1} \partial x^{2}}+\frac{1}{2 a} f_{2} \frac{\partial h}{\partial x^{1}}+\frac{1}{2 a} f_{1} \frac{\partial h}{\partial x^{2}}+\frac{1}{4 a^{2}} f_{1} f_{2} h\right] \partial_{2}^{y}
\end{align*}
$$

Theorem 1 Let $p \geq 2$. Then $\left(M, g_{\left(f_{1}, f_{2}, h\right)}\right)$ is $k$-Osserman for every admissible $k$.
Proof. Let be $X_{1}, X_{2}, X_{3}$ coordinate vector fields. By proposition $1, J\left(X_{1}\right) X_{3}=$ $R\left(X_{3}, X_{1}\right) X_{1}=0$ if $X_{1} \in \mathcal{Y}$. Thus $\mathcal{Y} \subset \operatorname{Ker}\left(J\left(X_{1}\right)\right)$. Furthermore, $\operatorname{range}\left(J\left(X_{2}\right)\right) \subset$ $\operatorname{span}\left\{R\left(\partial_{i}^{x}, \partial_{j}^{x}\right) \partial_{k}^{x}\right\} \subset \mathcal{Y}$. Thus $J\left(X_{1}\right) J\left(X_{2}\right)=0$.
If $\left\{X_{1}, X_{2}, \ldots, X_{k}\right\}$ is an orthonormal basis for $\pi \in G r_{(r, s)}\left(M, g_{\left(f_{1}, f_{2}, h\right)}\right)$, then we have

$$
J(\pi)^{2}=\sum_{i, j=1}^{k} g_{\left(f_{1}, f_{2}, h\right)}\left(X_{i}, X_{i}\right) g_{\left(f_{1}, f_{2}, h\right)}\left(X_{j}, X_{j}\right) J\left(X_{i}\right) J\left(X_{j}\right)=0
$$

Theorem 2 Let $p \geq 2$. The manifold $\left(\mathbf{R}^{4}, g_{\left(f_{1}, f_{2}, h\right)}\right)$ is a locally symmetric space if and only if the functions $f_{1}, f_{2}, h$ are solutions of the following partial differential equations in $\mathbf{R}^{2}$ :

$$
\begin{equation*}
\frac{\partial \Phi}{\partial x^{k}}+\frac{f_{k}}{2 a} \Phi=0, k=1,2 \tag{0.4}
\end{equation*}
$$

where we note

$$
\Phi\left(x^{1}, x^{2}\right)=\frac{1}{a}\left[\frac{\partial^{2} h}{\partial x^{1} \partial x^{2}}+\frac{1}{2 a} f_{2} \frac{\partial h}{\partial x^{1}}+\frac{1}{2 a} f_{1} \frac{\partial h}{\partial x^{2}}+\frac{1}{4 a^{2}} f_{1} f_{2} h\right]
$$

Proof. If we take in account this notation, we obtain by (0.3)

$$
R\left(\partial_{1}^{x}, \partial_{2}^{x}\right) \partial_{k}^{x}=(-1)^{k} \Phi\left(x^{1}, x^{2}\right) \partial_{3-k}^{y}, k=1,2
$$

Let $X_{k}=\alpha_{i}^{k} \partial_{i}^{x}, k=\overline{1,4}, i=\overline{1,4}$. The condition $\nabla_{X_{1}} R\left(X_{2}, X_{3}\right) X_{4}=0$ leads to

$$
\nabla_{\alpha_{i}^{1} \partial_{i}^{x}} R\left(\alpha_{j}^{2} \partial_{j}^{x}, \alpha_{l}^{3} \partial_{l}^{x}\right) \alpha_{s}^{4} \partial_{s}^{x}=0, i, j, k, s=\overline{1,4}
$$

Equivalently,

$$
\begin{gathered}
\alpha_{2}^{1} \alpha_{1}^{2} \alpha_{2}^{3} \alpha_{1}^{4} \nabla_{\partial_{2}^{x}} R\left(\partial_{1}^{x}, \partial_{2}^{x}\right) \partial_{1}^{x}+\alpha_{1}^{1} \alpha_{1}^{2} \alpha_{2}^{3} \alpha_{2}^{4} \nabla_{\partial_{1}^{x}} R\left(\partial_{1}^{x}, \partial_{2}^{x}\right) \partial_{2}^{x}+ \\
+\alpha_{2}^{1} \alpha_{2}^{2} \alpha_{1}^{3} \alpha_{1}^{4} \nabla_{\partial_{2}^{x}} R\left(\partial_{2}^{x}, \partial_{1}^{x}\right) \partial_{1}^{x}+\alpha_{1}^{1} \alpha_{2}^{2} \alpha_{1}^{3} \alpha_{2}^{4} \nabla_{\partial_{1}^{x}} R\left(\partial_{2}^{x}, \partial_{1}^{x}\right) \partial_{2}^{x}=0
\end{gathered}
$$

But

$$
\begin{gathered}
\nabla_{\partial_{1}^{x}} R\left(\partial_{1}^{x}, \partial_{2}^{x}\right) \partial_{2}^{x}=-\nabla_{\partial_{1}^{x}} R\left(\partial_{2}^{x}, \partial_{1}^{x}\right) \partial_{2}^{x}=\nabla_{\partial_{1}^{x}} \Phi \partial_{1}^{y}=\left(\frac{\partial \Phi}{\partial x^{1}}+\frac{f_{1}}{2 a} \Phi\right) \partial_{1}^{y} \\
\nabla_{\partial_{2}^{x}} R\left(\partial_{1}^{x}, \partial_{2}^{x}\right) \partial_{1}^{x}=-\nabla_{\partial_{2}^{x}} R\left(\partial_{2}^{x}, \partial_{1}^{x}\right) \partial_{1}^{x}=-\left(\frac{\partial \Phi}{\partial x^{2}}+\frac{f_{2}}{2 a} \Phi\right) \partial_{2}^{y}
\end{gathered}
$$

The proof is complete.
Corollary 1 If $h\left(x^{1}, x^{2}\right) \equiv C$ ( $h$ is a constant function), the conditions (0.4) for locally symmetry becames

$$
\left\{\begin{array}{l}
f_{1}^{\prime}\left(x^{1}\right) f_{2}\left(x^{2}\right)+\frac{1}{2 a} f_{1}^{2}\left(x^{1}\right) f_{2}\left(x^{2}\right)=0  \tag{0.5}\\
f_{2}^{\prime}\left(x^{2}\right) f_{1}\left(x^{1}\right)+\frac{1}{2 a} f_{2}^{2}\left(x^{2}\right) f_{1}\left(x^{1}\right)=0
\end{array}\right.
$$

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