# INVARIANT TRIPLE PRODUCTS 

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It is shown that the space of invariant trilinear forms on smooth representations of a semisimple Lie group is finite dimensional if the group is a product of hyperbolic groups. Explicit upper bounds are given which are attained in the case of induced representations. Applications to automorphic coefficients are given.

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## 1. Introduction

Let $G=\operatorname{PGL}_{2}(\mathbb{R})$ and let $\pi_{1}, \pi_{2}, \pi_{3}$ be irreducible admissible smooth representations of $G$. Then the space of $G$-invariant trilinear forms on $\pi_{1} \times \pi_{2} \times \pi_{3}$ is at most one dimensional. This has, in different contexts, been proved by Loke [9], Molčanov [10], and Oksak [11]. In this paper, we ask for such a uniqueness result in the context of arbitrary semisimple groups. We give evidence that for a given group $G$, uniqueness can only hold if $G$ is locally a product of hyperbolic groups. For such groups, we show uniqueness and for spherical vectors, we compute the invariant triple products explicitly.

By a conjecture of Jacquet's, which has been proved in [3], triple products on $\mathrm{GL}_{2}$ are related to special values of automorphic $L$-functions, see also [2, 4, 6, 7]. The conjecture/theorem says that the existence of nonzero triple products is equivalent to the nonvanishing of the corresponding triple $L$-function at the center of its functional equation.

The uniqueness of triple products in the $\mathrm{PGL}_{2}$-case mentioned above has been used in [1] to derive new bounds for automorphic $L^{2}$-coefficients. This can also be done for higher-dimensional hyperbolic groups, but, with the exception of the case treated in [1], the results do not exceed those in [8]. For completeness, we include these computations in the appendix.

## 2. Representations and integral formulae

Let $G$ be a connected semisimple Lie group with finite center. Fix a maximal compact subgroup $K$. Let $\widehat{G}$ and $\widehat{K}$ denote their unitary duals, that is, the sets of isomorphism
classes of irreducible unitary representations of $G$, respectively, $K$. Let $\pi$ be a continuous representation of $G$ on a locally convex topological vector space $V_{\pi}$. Let $V_{\pi}^{\prime}$ be the space of all continuous linear forms on $V_{\pi}$ and let $V_{\pi}^{\infty}$ be the space of smooth vectors, that is,

$$
\begin{equation*}
V_{\pi}^{\infty}=\left\{v \in V_{\pi}: x \longmapsto \alpha(\pi(x) v) \text { is smooth } \forall \alpha \in V_{\pi}^{\prime}\right\} \tag{2.1}
\end{equation*}
$$

The representation $\pi$ is called smooth if $V_{\pi}=V_{\pi}^{\infty}$.
A representation $\left(\pi, V_{\pi}\right)$ is called admissible if for each $\tau \in \hat{K}$, the space $\operatorname{Hom}_{K}\left(V_{\tau}, V_{\pi}\right)$ is finite-dimensional. Let $\widehat{G}_{\text {adm }}$ be the admissible dual, that is, the set of infinitesimal isomorphism classes of irreducible admissible representations. A representation $\pi$ in $\widehat{G}_{\text {adm }}$ is called a class one or spherical representation if it contains $K$-invariant vectors. In that case, the space $V_{\pi}^{K}$ of $K$-invariant vectors is one-dimensional. This is trivial for principal series representations (see below) and follows generally from Casselman's subrepresentation theorem which says that every $\pi \in \widehat{G}_{\text {adm }}$ is equivalent to a subrepresentation of a principal series representation.

The Iwasawa decomposition $G=A N K$ gives smooth maps

$$
\begin{align*}
& \underline{a}: G \longrightarrow A, \\
& \underline{n}: G \longrightarrow N,  \tag{2.2}\\
& \underline{k}: G \longrightarrow K,
\end{align*}
$$

such that for every $x \in G$, one has $x=\underline{a}(x) \underline{n}(x) \underline{k}(x)$. As an abbreviation, we also define $\underline{a n}(x)=\underline{a}(x) \underline{n}(x)$. Let $\mathfrak{g}_{\mathbb{R}}, \mathfrak{a}_{\mathbb{R}}, \mathfrak{n}_{\mathbb{R}}, \mathfrak{k}_{\mathbb{R}}$ denote the Lie algebras of $G, A, N, K$ and let $\mathfrak{g}, \mathfrak{a}, \mathfrak{n}$, $\mathfrak{k}$ be their complexifications.

For $x \in G$ and $k \in K$, we define

$$
\begin{equation*}
k^{x} \stackrel{\operatorname{def}}{=} \underline{k}(k x) \tag{2.3}
\end{equation*}
$$

Lemma 2.1. The rule $k \mapsto k^{x}$ defines a smooth (right) group action of $G$ on $K$.
Proof. The map $\underline{k}$ gives a diffeomorphism $A N \backslash G \rightarrow K$ and the action under consideration is just the natural right action of $G$ on $A N \backslash G$.

Lemma 2.2. For $f \in C(K)$ and $y \in G$, one has the integral formula

$$
\begin{equation*}
\int_{K} f(k) d k=\int_{K} \underline{a}(k y)^{2 \rho} f\left(k^{y}\right) d k \tag{2.4}
\end{equation*}
$$

or

$$
\begin{equation*}
\int_{K} f\left(k^{y}\right) d k=\int_{K} \underline{a}\left(k y^{-1}\right)^{2 \rho} f(k) d k \tag{2.5}
\end{equation*}
$$

Here $\rho \in \mathfrak{a}^{*}$ is the modular shift, that is, $a^{2 \rho}=\operatorname{det}(a \mid \mathfrak{n})$.
Proof. The Iwasawa integral formula implies that

$$
\begin{equation*}
\int_{G} g(x) d x=\int_{A N} \int_{K} g(a n k) d k d a n . \tag{2.6}
\end{equation*}
$$

Let $f \in C(K)$ and choose a function $\eta \in C_{c}(A N)$ such that $\eta \geq 0$ and $\int_{A N} \eta(a n) d a d n=1$. Let $g(x)=\eta(\underline{a n}(x)) f(\underline{k}(x))$. Then $g \in C_{c}(G)$ and

$$
\begin{equation*}
\int_{G} g(x) d x=\int_{A N} \eta(a n) d a n \int_{K} f(k) d k=\int_{K} f(k) d k \tag{2.7}
\end{equation*}
$$

On the other hand, since $G$ is unimodular, this also equals

$$
\begin{align*}
\int_{G} g(x y) d x & =\int_{G} \eta(\underline{a n}(x y)) f(\underline{k}(x y)) d x \\
& =\int_{A N} \int_{K} \eta(\underline{a n}(\operatorname{anky})) f(\underline{k}(k y)) d a n k \\
& =\int_{K}\left(\int_{A N} \eta(\underline{a n}(\operatorname{ank} y)) d a n\right) f\left(k^{y}\right) d k \\
& =\int_{K}\left(\int_{A N} \eta(\operatorname{anan}(k y)) d a n\right) f\left(k^{y}\right) d k  \tag{2.8}\\
& =\int_{K} \underline{a}(k y)^{2 \rho}\left(\int_{A N} \eta(a n) d a n\right) f\left(k^{y}\right) d k \\
& =\int_{K} \underline{a}(k y)^{2 \rho} f\left(k^{y}\right) d k .
\end{align*}
$$

The second assertion follows from the first by replacing $f$ with $\tilde{f}(k)=f\left(k^{y}\right)$ and then $y$ with $y^{-1}$.

Let $M$ be the centralizer of $A$ intersected with $K$, then $P=M A N$ is a minimal parabolic subgroup of $G$. The inclusion map $K \hookrightarrow G$ induces a diffeomorphism $M \backslash K \rightarrow P \backslash G$ and in this way, we get a smooth $G$-action on $M \backslash K$. An inspection shows that this action is given by $M k \mapsto M k^{x}$ for $k \in K, x \in G$.

## 3. Trilinear products

Let $\pi_{1}, \pi_{2}, \pi_{3}$ be three admissible smooth representations of the group $G$ and let $\mathscr{T}$ : $V_{\pi_{1}} \times V_{\pi_{2}} \times V_{\pi_{3}} \rightarrow \mathbb{C}$ be a continuous $G$-invariant trilinear form, that is,

$$
\begin{equation*}
\mathscr{T}\left(\pi_{1}(x) v_{1}, \pi_{2}(x) v_{2}, \pi_{3}(x) v_{3}\right)=\mathscr{T}\left(v_{1}, v_{2}, v_{3}\right) \tag{3.1}
\end{equation*}
$$

for all $v_{j} \in V_{\pi_{j}}$ and every $x \in G$.
We want to understand the space of all trilinear forms $\mathscr{T}$ as above. In this paper, we will only consider principal series representations, the general case will be considered later. So we assume that $\pi_{1}, \pi_{2}, \pi_{3}$ are principal series representations. This means that there are given a minimal parabolic $P=M A N$, irredicible representations $\sigma_{j} \in \widehat{M}$, and $\lambda_{j} \in \mathfrak{a}^{*}$ for $j=1,2,3$. Each pair $\left(\sigma_{j}, \lambda_{j}\right)$ induces a continuous group homomorphism $P \rightarrow \operatorname{GL}\left(V_{\sigma_{j}}\right)$ by man $\mapsto a^{\lambda_{j}+\rho} \sigma(m)$, which in turn defines a $G$-homogeneous vector bundle $E_{\sigma_{j}, \lambda_{j}}$ over $P \backslash G$. The representation $\pi_{j}$ is the $G$-representation on the space of smooth sections $\Gamma^{\infty}\left(E_{\sigma_{j}, \lambda_{j}}\right)$ of that bundle. In other words, $\pi_{j}$ lives on the space of all $C^{\infty}$ functions $f: G \rightarrow V_{\sigma_{j}}$ with

$$
\begin{equation*}
f(\operatorname{man} x)=a^{\lambda_{j}+\rho} \sigma_{j}(m) f(x) \tag{3.2}
\end{equation*}
$$

for all $m \in M, a \in A, n \in N, x \in G$. The representation $\pi_{j}$ is defined by $\pi_{j}(y) f(x)=$ $f(x y)$. Every such $f$ is uniquely determined by its restriction to $K$ which satisfies $f(m k)=$ $\sigma_{j}(m) f(k)$, that is, $f$ is a section of the $K$-homogeneous bundle $E_{\sigma_{j}}$ on $M \backslash K$ induced by $\sigma_{j}$. So the representation space can be identified with $V_{\pi_{j}} \cong \Gamma^{\infty}\left(E_{\sigma_{j}}\right)$. Thus a trilinear form $\mathscr{T}$ is a distribution on the vector bundle $E_{\sigma}=E_{\sigma_{1}} \boxtimes E_{\sigma_{2}} \boxtimes E_{\sigma_{3}}$ over $M \backslash K \times M \backslash K \times M \backslash K$. Here $\boxtimes$ denotes the outer tensor product.

For $f_{1}, f_{2}, f_{3} \in C^{\infty}(M \backslash K)$, we write $\mathscr{T}\left(f_{1}, f_{2}, f_{3}\right)$ for the expression

$$
\begin{equation*}
\int_{M \backslash K \times M \backslash K \times M \backslash K} \phi\left(k_{1}, k_{2}, k_{3}\right)\left[f_{1}\left(k_{1}\right) \boxtimes f_{2}\left(k_{2}\right) \boxtimes f_{3}\left(k_{3}\right)\right] d k_{1} d k_{2} d k_{3}, \tag{3.3}
\end{equation*}
$$

where $\phi$ is the kernel of $\mathscr{T}$.
The group $G$ is called a real hyperbolic group if it is locally isomorphic to $\operatorname{SO}(d, 1)$ for some $d \geq 2$.

On $Y=(P \backslash G)^{3}$, we consider the $G^{3}$-homogeneous vector bundle $E_{\sigma, \lambda}$ given by $E_{\sigma, \lambda}=$ $E_{\sigma_{1}, \lambda_{1}} \boxtimes E_{\sigma_{2}, \lambda_{2}} \boxtimes E_{\sigma_{3}, \lambda_{3}}$. Next $Y$ can be viewed as a $G$-space via the diagonal action and so $E_{\sigma, \lambda}$ becomes a $G$-homogeneous line bundle on $Y$.

We are going to impose the following condition on the induction parameters $\lambda_{1}, \lambda_{2}$, $\lambda_{3}$. We assume that $\sum_{j=1}^{3} \varepsilon_{j}\left(\lambda_{j}+\rho\right) \neq 0$ for any choice of $\varepsilon_{j} \in\{ \pm 1\}$. In other words, this means that
(i) $\lambda_{1}+\lambda_{2}+\lambda_{3}+3 \rho \neq 0$,
(ii) $\lambda_{1}+\lambda_{2}-\lambda_{3}+\rho \neq 0$,
(iii) $\lambda_{1}-\lambda_{2}-\lambda_{3}-\rho \neq 0$.

Theorem 3.1. Assume the parameters $\lambda_{1}, \lambda_{2}, \lambda_{3}$ satisfy the above condition. Let $Y$ be the $G$-space $(P \backslash G)^{3}$. If there is an open $G$-orbit in $Y$, then the dimension of the space of invariant trilinear forms on smooth principal series representations is less than or equal to

$$
\begin{equation*}
\sum_{o} \operatorname{dim}\left(\sigma_{1} \otimes \sigma_{2} \otimes \sigma_{3}\right)^{M_{o}} \tag{3.4}
\end{equation*}
$$

where the sum runs over all open orbits o and $M_{o}$ is the stabilizer group of a point in the orbit o which is chosen so that $M_{o}$ is a subgroup of $M$. If all induction parameters are imaginary (unitary induction), then there is equality.

In particular, if $\pi_{1}, \pi_{2}, \pi_{3}$ are class-one representations, then the dimension is less than or equal to the number of open orbits in $Y$.

There is an open orbit if and only if $G$ is locally isomorphic to a product of hyperbolic groups.

For $G=\mathrm{SO}(2,1)^{0}$, the number of open orbits is 2 , for $G=\mathrm{SO}(2,1)$ or $G=\mathrm{SO}(d, 1)^{0}$, $d \geq 3$, the number is 1 . Here $\mathrm{SO}(d, 1)^{0}$ is the connected component of the Lie group $\mathrm{SO}(d, 1)$.

The Proof is based on the following lemma.
Lemma 3.2. Let $G$ be a Lie group and $H$ a closed subgroup. Let $X=G / H$ and let $E \rightarrow X$ be a smooth $G$-homogeneous vector bundle. Let $\mathscr{T}$ be a distribution on $E$, that is, a continuous linear form on $\Gamma_{c}^{\infty}(E)$. Suppose that $\mathcal{T}$ is $G$-invariant, that is, $\mathcal{T}(g \cdot s)=\mathscr{T}(s)$ for every $s \in \Gamma_{c}^{\infty}(E)$. Then $\mathscr{T}$ is given by a smooth $G$-invariant section of the dual bundle $E^{*}$.

Let $\left(\sigma, V_{\sigma}\right)$ be the representation of $H$ on the fibre $E_{e H}$ and let $\left(\sigma^{*}, V_{\sigma^{*}}\right)$ be its dual. Then the space of all $G$-invariant distributions on $E$ has dimension equal to the dimension $\operatorname{dim} V_{\sigma^{*}}^{H}$ of $H$-invariants. So in particular, if $\sigma$ is irreducible, this dimension is zero unless $\sigma$ is trivial, in which case the dimension is one.

Proof. Equation $\mathscr{T}(g \cdot s)=\mathscr{T}(s)$, that is, $g \cdot \mathscr{T}=\mathscr{T}$ for all $g \in G$ implies that $X \cdot \mathscr{T}=0$ for every $X \in \mathfrak{g}_{\mathbb{R}}$, the real Lie algebra of $G$. Let $\mathfrak{h}_{\mathbb{R}}$ be the Lie algebra of $H$ and choose a complementary space $\mathfrak{p}_{\mathbb{R}}$ for $\mathfrak{h}_{\mathbb{R}}$ such that $\mathfrak{g}_{\mathbb{R}}=\mathfrak{h}_{\mathbb{R}} \oplus \mathfrak{p}_{\mathbb{R}}$. Let $X_{1}, \ldots, X_{n}$ be a basis of $V$ and let

$$
\begin{equation*}
D=X_{1}^{2}+X_{2}^{2}+\cdots+X_{n}^{2} \in \mathrm{U}\left(\mathfrak{g}_{\mathbb{R}}\right) . \tag{3.5}
\end{equation*}
$$

We show that $D$ induces an elliptic differential operator on $E$. By $G$-homogeneity, it suffices to show this at a single point. So let $P=\exp \left(\mathfrak{p}_{\mathbb{R}}\right)$. In a neighborhood $U$ of the unit in $G$, there are smooth maps $\underline{h}: U \rightarrow H$ and $\underline{p}: U \rightarrow P$ such that $\underline{h}(x) \underline{p}(x)=x$ for $x \in U$.

The sections of $E$ can be identified with the smooth maps $s: G \rightarrow V_{\sigma}$ with $s(h x)=$ $\sigma(h) s(x)$ for $h \in H$ and $x \in G$. We can attach to each section $s$ a map $f_{s}$ on $\mathfrak{p}$ with values in $V_{\sigma}$ by $f_{s}(Y)=s(\exp (Y))$. The action of $X \in \mathfrak{p}_{\mathbb{R}}$ on the section $s$ is described by

$$
\begin{align*}
f_{X s}(Y) & =\left.\frac{d}{d t}\right|_{t=0} s(\exp (Y) \exp (t X))  \tag{3.6}\\
& =\left.\frac{d}{d t}\right|_{t=0} \sigma(\underline{h}(\exp (Y) \exp (t X)) s(\underline{p}(\exp (Y) \exp ((t X)) .
\end{align*}
$$

Let $A_{X}(Y)=d /\left.d t\right|_{t=0} \sigma\left(\underline{h}(\exp (Y) \exp (t X)) \in \operatorname{End}\left(V_{\sigma}\right)\right.$. Then the Leibniz rule implies that

$$
\begin{equation*}
f_{X s}(Y)=A_{X}(Y) f_{s}(Y)+X f_{s}(Y) \tag{3.7}
\end{equation*}
$$

The first summand is of order zero and the second is of order one. Moreover, the second summand at $Y=0$ coincides with the coordinate-derivative in the direction of $X$. This implies that the leading symbol of $D$ at $e H$ is $\xi_{1}^{2}+\cdots+\xi_{n}^{2}$ and so $D$ is elliptic. The distributional equation $D \mathscr{T}=0$ then implies that $\mathscr{T}$ is given by a smooth section.

For the second assertion of the lemma, recall that a $G$-invariant section is uniquely determined by its restriction to the point eH which must be invariant under $H$.

For the proof of the theorem, we will need to investigate the $G$-orbit structure of $Y=(P \backslash G)^{3}$. First note that since the map $M k \mapsto P k$ is a $K$-isomorphism from $M \backslash K$ to $P \backslash G$, the $K$-orbit of every $y \in Y$ contains an element of the form $\left(y_{1}, y_{2}, 1\right)$. Hence the $P=M A N$-orbit structure of $(P \backslash G)^{2}$ is the same as the $G$-orbit structure of $Y$. By the Bruhat decomposition, the $P$-orbits in $P \backslash G$ are parametrized by the Weyl group $W=W(G, A)$, where the unique open orbit is given by $P w_{0} P$, here $w_{0}$ is the long element of the Weyl group. Note that the $P$-stabilizer of $P w_{0} \in P \backslash G$ equals $A M$. This implies that the $G$-orbits in $Y$ of maximal dimension are in bijection to the $A M$-orbits in
$P \backslash G$ of maximal dimension via the map $P x A M \mapsto\left(x, w_{0}, 1\right) \cdot G$. Again by Bruhat decomposition, it follows that the latter are contained in the open cell $P w_{0} P=P w_{0} N$. So the $G$-orbits of maximal dimension in $Y$ are in bijection to the $A M$-orbits in $N$ of maximal dimension, where $A M$ acts via the adjoint action. The exponential map exp : $\mathfrak{n}_{\mathbb{R}} \rightarrow N$ is an $A M$-equivariant bijection, so we are finally looking for the $A M$-orbit structure of the linear adjoint action on $\mathfrak{n}_{\mathbb{R}}$.

We will now prove that there is an open orbit if and only if $G$ is locally a product of real hyperbolic groups. So suppose that $Y$ contains an open orbit. Then $\mathfrak{n}_{\mathbb{R}}$ contains an open $A M$-orbit, say $A M \cdot X_{0}$. Let $\phi^{+}$be the set of all positive restricted roots on $\mathfrak{a}=\operatorname{Lie}(A)$. Decompose $\mathfrak{n}_{\mathbb{R}}$ into the root spaces

$$
\begin{equation*}
\mathfrak{n}_{\mathbb{R}}=\bigoplus_{\alpha \in \phi^{+}} \mathfrak{n}_{\mathbb{R}, \alpha} \tag{3.8}
\end{equation*}
$$

On each $\mathfrak{n}_{\mathbb{R}, \alpha}$, install an $M$-invariant norm $\|\cdot\|_{\alpha}$. This is possible since $M$ is compact. Consider the map

$$
\begin{align*}
\psi & : \mathfrak{n}_{\mathbb{R}} \longrightarrow \prod_{\alpha \in \phi^{+}} \mathbb{R}  \tag{3.9}\\
x & \longmapsto \prod_{\alpha}\|x\|_{\alpha}
\end{align*}
$$

Since the orbit $A M \cdot X_{0}$ is open, the image $\psi\left(A M \cdot X_{0}\right)$ of the orbit must contain a nonempty open set. Away from the set $\left\{X \in \mathfrak{n}_{\mathbb{R}}: \exists \alpha:\|x\|_{\alpha}=0\right\}$, the map $\psi$ can be chosen differentiable. Since the norms are invariant under $M$, one gets a smooth map

$$
\begin{gather*}
A \longrightarrow \mathbb{R}^{\left|\phi^{+}\right|} \\
a \longmapsto \psi\left(a \cdot X_{0}\right), \tag{3.10}
\end{gather*}
$$

whose image contains an open set. This can only happen if the dimension of $A$ is at least as big as $\left|\phi^{+}\right|$and the latter implies that $G$ is locally a product of real rank-one groups. Now by Araki's table (see [5, pages 532-534]), one knows that these real rank-one groups must all be hyperbolic, because otherwise there would be two different root lengths.

For the converse direction, let $G$ be locally isomorphic to $\mathrm{SO}(d, 1)$. We have to show that there is an open $A M$-orbit in $\mathfrak{n}_{\mathbb{R}}$. This, however, is clear as the action of $A M$ on $\mathfrak{n}_{\mathbb{R}}$ is the natural action of $\mathbb{R}_{+}^{\times} \times \operatorname{SO}(d-1)$ on $\mathbb{R}^{d-1}$, hence there are two orbits, the zero orbit and one open orbit.

We will now show that if there is an open orbit, then there are no invariant distributions supported on lower-dimensional orbits. For this, it suffices to consider the case $G=\mathrm{SO}^{0}(d, 1)$. For simplicity, we only consider the trivial bundle, that is, functions instead of sections. The general case is similar. The orbit structure is as follows. For $d=2$, the group $M=\operatorname{SO}(d-1)$ is trivial, and so there are two open $A M$-orbits in $N \cong \mathbb{R}$ in this case. If $d>2$, there is only one open orbit. Since the proof is very similar in the case $d=2$, we will now restrict ourselves to the case $d>2$. The open orbit [ $\left.w_{0} n_{0}, w_{0}, 1\right]$, given by some $n_{0} \in N$, contains in its closure the orbits [ $w_{0}, w_{0}, 1$ ] and [ $w_{0}, 1,1$ ] which in turn
contain in their closures the orbit $[1,1,1]$,

$$
[1,1,1]
$$

Let $\mathbb{C}=\left[w_{0} n_{0}, w_{0}, 1\right]$ be the open orbit. Let $x_{1}=\left(w_{0}, w_{0}, 1\right)$, then one has $\left[w_{0}, w_{0}, 1\right]=$ $x_{1} \cdot G$. Consider $A \cong \mathbb{R}_{+}^{\times}$as a subset of $\mathbb{R}$ suitably normalized, then one can write

$$
\begin{equation*}
x_{1}=\lim _{a \rightarrow 0} x_{0} \cdot a m . \tag{3.12}
\end{equation*}
$$

Let $\mathscr{T}$ be a $G$-invariant distribution supported on the closure of the orbit of $x_{1}$. Since $\mathscr{T}$ is $G$-invariant, it satisfies $X \cdot \mathscr{T}=0$ for every $X \in \mathfrak{g}$. Hence the wave front set $W F(\mathscr{T}) \subset$ $T^{*} Y$ is a $G$-invariant subset of the normal bundle of the manifold $x_{1} \cdot G$. This implies that $\mathscr{T}$ is of order zero along the manifold $x_{1} \cdot G$. By the $G$-invariance, it follows that $\mathscr{T}$ is of the form

$$
\begin{equation*}
\mathscr{T}(f)=\int_{x_{1} \cdot G} D\left(^{x} f\right)\left(x_{1}\right) d x+R \tag{3.13}
\end{equation*}
$$

where $R$ is supported in $[1,1,1]$. Further, $D$ is a differential which we can assume to be $G$-equivariant. Then $D(f)\left(x_{1}\right)$ is of the form

$$
\begin{equation*}
\left.D_{1}(m) f\left(x_{0} \cdot a m\right)\right|_{a=0} \tag{3.14}
\end{equation*}
$$

where $D_{1}(m)$ is a differential operator in the variable $a$. Since $D$ is $G$-equivariant, we may replace $f$ with ${ }^{a_{0}} f$ for some $a_{0} \in A$. Since $x_{1} a_{0}=x_{1}$, we get that the above is the same as

$$
\begin{equation*}
\left.D_{1}(m) f\left(x_{0} \cdot a a_{0} m\right)\right|_{a=0} . \tag{3.15}
\end{equation*}
$$

This implies that $D_{1}$ must be of order zero and so $\mathscr{T}$ is of order zero. Restricted to the orbit $x_{1} \cdot G \cong A M \backslash G$, the distribution $\mathscr{T}$ is given by an integral of the form $\int_{A M \backslash G} \phi(y) f\left(x_{1} y\right) d y$. (Note that we use the notation without the dot again.) Invariance implies that $\phi$ is constant. If $\mathscr{T}$ is nonzero, then $y \mapsto f\left(x_{1} y\right)$ must be left invariant under $A M$, which implies that $\lambda_{1}+\lambda_{2}-\lambda_{3}+\rho=0$, a case we have excluded. So $\mathscr{T}$ must be zero. This shows that any invariant distribution which is zero on the open orbit also vanishes on $x_{1} \cdot G$. The remaining orbits are dealt with in a similar fashion. To prove the theorem, it remains to show the existence of invariant distributions in the case of unitary parameters. For this, we change our point of view and consider sections of $L_{\lambda}$ no longer as functions on $Y$, but as functions on $G^{3}$ with values in $V_{\sigma}=V_{\sigma_{1}} \otimes V_{\sigma_{2}} \otimes V_{\sigma_{3}}$ which spit out $a^{\lambda+\rho} \sigma(m)$ on the left. We induce this in the notation by writing $f\left(x_{0} y\right)$ instead of $f\left(x_{0} \cdot y\right)$. On a given orbit of maximal dimension, there is a standard invariant distribution which, by the lemma,
is unique up to scalars and is given by

$$
\begin{equation*}
\mathscr{T}_{\text {st }} f=\alpha\left(\int_{G} f\left(x_{0} y\right) d y\right) \tag{3.16}
\end{equation*}
$$

where $\alpha$ is a linear functional on the space of $M_{o}$-invariants. In order to show that this extends to a distribution on $Y$, we need to show that the defining integral converges for all $f \in \Gamma^{\infty}\left(E_{\sigma, \lambda}\right)$. This integral equals

$$
\begin{align*}
\int_{G} f\left(x_{0} y\right) d y & =\int_{G} f\left(w_{0} n_{0} y, w_{0} y, y\right) d y \\
& =\int_{G} \underline{a}\left(w_{0} n_{0} y\right)^{\rho+\lambda_{1}} \underline{a}\left(w_{0} y\right)^{\rho+\lambda_{2}} \underline{a}(y)^{\rho+\lambda_{3}} f\left(\underline{k}\left(x_{0} y\right)\right) d y . \tag{3.17}
\end{align*}
$$

Since $f$ is bounded on $K^{3}$, it suffices to show the following lemma.
Lemma 3.3. Let $G=\operatorname{SO}(d, 1)^{0}$ and let $k_{0}=\underline{k}\left(w_{0} n_{0}\right)$. Then

$$
\begin{equation*}
\int_{G} \underline{a}\left(w_{0} n_{0} y\right)^{\rho} \underline{a}\left(w_{0} y\right)^{\rho} \underline{a}(y)^{\rho} d y<\infty . \tag{3.18}
\end{equation*}
$$

Conjecture 3.4. The assertion of the lemma should hold for any semisimple group $G$ with finite center and $n_{0} \in N$ generic.

The conjecture would imply that if $G$ is not locally a product of hyperbolic groups, then the space of invariant trilinear forms on principal series representations is infinitedimensional.

Proof of Lemma 3.3. Replace the integral over $G$ by an integral over ANK using the Iwasawa decomposition. Since $\underline{a}(x k)=\underline{a}(x)$ for $x \in G$ and $k \in K$, the $K$-factor is irrelevant and we have to show that

$$
\begin{equation*}
\int_{A N} \underline{a}\left(k_{0} a n\right)^{\rho} \underline{a}\left(w_{0} a n\right)^{\rho} a^{\rho} d a d n<\infty . \tag{3.19}
\end{equation*}
$$

Now write $w_{0} n_{0}=a^{\prime} n^{\prime} k_{0}$. Then $k_{0} a n=\left(a^{\prime} n^{\prime}\right)^{-1} w_{0} n_{0} a n$, and therefore $\underline{a}\left(k_{0} a n\right)=$ $\left(a^{\prime}\right)^{-1} \underline{a}\left(w_{0} n_{0} a n\right)$, so it suffices to show that

$$
\begin{equation*}
\int_{A N} \underline{a}\left(w_{0} n_{0} a n\right)^{\rho} \underline{a}\left(w_{0} a n\right)^{\rho} a^{\rho} d a d n<\infty . \tag{3.20}
\end{equation*}
$$

Next note that $w_{0} a=a^{-1} w_{0}$ and so we have $\underline{a}\left(w_{0} a n\right)^{\rho}=a^{-\rho} \underline{a}\left(w_{0} n\right)^{\rho}$ as well as $\underline{a}\left(w_{0} n_{0} a n\right)^{\rho}=$ $a^{-\rho} \underline{a}\left(w_{0} n_{0}^{a} n\right)^{\rho}$, where $n_{0}^{a}=a^{-1} n_{0} a$. We need to show that

$$
\begin{equation*}
\int_{A N} \underline{a}\left(w_{0} n_{0}^{a} n\right)^{\rho} \underline{a}\left(w_{0} n\right)^{\rho} a^{-\rho} d a d n<\infty \tag{3.21}
\end{equation*}
$$

This is the point where we have to make things more concrete. Let $J$ be the diagonal $(d+1) \times(d+1)$-matrix with diagonal entries $(1, \ldots, 1,-1)$. Then $\mathrm{SO}(d, 1)$ is the group of
real matrices $g$ with $g^{t} J g=J$. Writing $g=\left(\begin{array}{cc}A & b \\ c & d\end{array}\right)$ with $A \in \operatorname{Mat}_{d}(\mathbb{R})$, this amounts to

$$
\begin{gather*}
A^{t} A-c^{t} c=1, \\
A^{t} b-c^{t} d=0,  \tag{3.22}\\
d^{2}-b^{t} b=1
\end{gather*}
$$

The connected component $\operatorname{SO}(d, 1)^{0}$ consists of all matrices $g$ as above with $d>0$. The maximal compact subgroup $K$ can be chosen to be

$$
\left(\begin{array}{c|c}
\mathrm{SO}(d) &  \tag{3.23}\\
\hline & 1
\end{array}\right)
$$

and $M$ as

$$
\left(\begin{array}{cc|c}
\mathrm{SO}(d-1) & & 0  \tag{3.24}\\
& 1 & 0 \\
\hline 0 & 0 & 1
\end{array}\right)
$$

Further, we can choose $A$ and $N$ as follows:

$$
\begin{gather*}
A=\left\{\left(\begin{array}{cc|c}
1 & & \\
& \alpha & \beta \\
\hline & \beta & \alpha
\end{array}\right): \alpha>0, \alpha^{2}-\beta^{2}=1\right\}, \\
N=\left\{n(x)=\left(\begin{array}{cc|c}
1 & -x & x \\
x^{t} & 1-\frac{|x|^{2}}{2} & D \frac{|x|^{2}}{2} \\
\hline x^{t} & -\frac{|x|^{2}}{2} & 1+\frac{|x|^{2}}{2}
\end{array}\right): x \in \mathbb{R}^{d-1}\right\}, \tag{3.25}
\end{gather*}
$$

where we have written $|x|^{2}=x_{1}^{2}+x_{2}^{2}+\cdots+x_{d-1}^{2}$. Note that the Lie algebra of $A$ is generated by $H=\left(\begin{array}{cc}1 & 0 \\ 0 & 1\end{array}\right)$. One derives an explicit formula for the ANK-decomposition. In particular, if $g=\left(\begin{array}{cc}A & b \\ c & d\end{array}\right)$ and

$$
\underline{a}(g)=\left(\begin{array}{ll|l}
1 & &  \tag{3.26}\\
& \alpha & \beta \\
\hline & \beta & \alpha
\end{array}\right),
$$

then

$$
\begin{equation*}
\underline{a}(g)^{\rho}=(\alpha+\beta)^{(d-1) / 2}=\left(\frac{\left(d+b_{d}\right)}{1+b_{1}^{2}+\cdots+b_{d-1}^{2}}\right)^{(d-1) / 2} . \tag{3.27}
\end{equation*}
$$

The Weyl representative element can be chosen to be

$$
w_{0}=\left(\begin{array}{ccccc|c}
1 & & & & &  \tag{3.28}\\
& \ddots & & & & \\
& & 1 & & & \\
& & & -1 & & \\
& & & & -1 & \\
\hline & & & & & 1
\end{array}\right) .
$$

So that with $n=n(x)$ for $x \in \mathbb{R}^{d-1}$, we have

$$
\begin{align*}
\underline{a}\left(w_{0} n\right)^{\rho} & =\underline{a}\left(\begin{array}{cc|c}
* & * & x_{1} \\
* & * & \vdots \\
* & * & x_{d-2} \\
* & * & -x_{d-1} \\
* & * & -\frac{|x|^{2}}{2} \\
\hline * & * & 1+\frac{|x|^{2}}{2}
\end{array}\right)^{\rho}  \tag{3.29}\\
& =\left(\frac{1}{1+x_{1}^{2}+\cdots+x_{d-1}^{2}}\right)^{(d-1) / 2} .
\end{align*}
$$

We choose $n_{0}=n(1,0, \ldots, 0)$ and get with $a=\exp (t H)$ that

$$
\begin{equation*}
\underline{a}\left(w_{0} n_{0}^{a} n\right)^{\rho}=\left(\frac{1}{1+\left(e^{-t}+x_{1}\right)^{2}+x_{2}^{2}+\cdots+x_{d-1}^{2}}\right)^{(d-1) / 2} . \tag{3.30}
\end{equation*}
$$

With $n=d-1$ and $a^{-\rho}=e^{-(n / 2) t}$, our assertion boils down to

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} \int_{\mathbb{R}} e^{-(n / 2) t}\left(1+|x|^{2}\right)^{-n / 2}\left(1+\left|e^{-t} v_{1}+x\right|^{2}\right)^{-n / 2} d t d x<\infty . \tag{3.31}
\end{equation*}
$$

Consider first the case $n=1$ and the integral over $x<0$ :

$$
\begin{align*}
\int_{\mathbb{R}} \int_{-\infty}^{0} & e^{-(1 / 2) t}\left(1+x^{2}\right)^{-1 / 2}\left(1+\left(e^{-t}+x\right)^{2}\right)^{-1 / 2} d x d t \\
& =\int_{\mathbb{R}} \int_{0}^{\infty} e^{-(1 / 2) t}\left(1+x^{2}\right)^{-1 / 2}\left(1+\left(x-e^{-t}\right)^{2}\right)^{-1 / 2} d x d t  \tag{3.32}\\
& =\int_{\mathbb{R}} \int_{-e^{-t}}^{\infty} e^{-(1 / 2) t}\left(1+\left(x+e^{-t}\right)^{2}\right)^{-1 / 2}\left(1+(x)^{2}\right)^{-1 / 2} d x d t .
\end{align*}
$$

Thus it suffices to show the convergence of

$$
\begin{equation*}
\int_{\mathbb{R}} \int_{0}^{\infty} e^{-(1 / 2) t}\left(1+x^{2}\right)^{-1 / 2}\left(1+\left(x+e^{-t}\right)^{2}\right)^{-1 / 2} d t d x \tag{3.33}
\end{equation*}
$$

Setting $y=e^{-t}$, we see that this integral equals

$$
\begin{align*}
& \iint_{0}^{\infty} y^{-1 / 2}\left(1+x^{2}\right)^{-1 / 2}\left(1+(x+y)^{2}\right)^{-1 / 2} d y d x \\
& \quad=\int_{0}^{\infty} \int_{x}^{\infty}(y-x)^{-1 / 2}\left(1+x^{2}\right)^{-1 / 2}\left(1+y^{2}\right)^{-1 / 2} d y d x \tag{3.34}
\end{align*}
$$

Since the mapping

$$
\begin{equation*}
x \longmapsto \int_{x}^{\infty}(y-x)^{-1 / 2}\left(1+x^{2}\right)^{-1 / 2}\left(1+y^{2}\right)^{-1 / 2} d y \tag{3.35}
\end{equation*}
$$

is continuous, the integral over $0<x<1$ converges. It remains to show the convergence of

$$
\begin{align*}
& \int_{1}^{\infty} \int_{x}^{\infty}(y-x)^{-1 / 2}\left(1+x^{2}\right)^{-1 / 2}\left(1+y^{2}\right)^{-1 / 2} d y d x \\
& \quad=\iint_{1}^{\infty} x^{1 / 2}(v-1)^{-1 / 2}\left(1+x^{2}\right)^{-1 / 2}\left(1+v^{2} x^{2}\right)^{-1 / 2} d v d x  \tag{3.36}\\
& \quad \leq 3^{-1 / 2} \iint_{1}^{\infty} x^{1 / 2}(v-1)^{-1 / 2}\left(1+x^{2}\right)^{-1}\left(1+v^{2}\right)^{-1 / 2} d v d x<\infty .
\end{align*}
$$

Here we have used the substitution $y=v x$ and the fact that for $a, b \geq 1$, one has ( $1+$ a) $(1+b) \leq 3(1+a b)$.

Now for the case $n>1$, using polar coordinates, we compute

$$
\begin{align*}
& \left.\left.\int_{\mathbb{R}^{n-1}} \quad \iint_{\mathbb{R}} e^{-t n / 2}\left(1+x^{2}+\left|x_{r}\right|^{2}\right)^{-n / 2}(1+) e^{-t}+x\right)^{2}+\left|x_{r}\right|^{2}\right)^{-n / 2} d t d x d x_{r} \\
& \left.\left.\quad=C \int_{0}^{\infty} \iint_{\mathbb{R}} e^{-t n / 2} r^{n-2}\left(1+x^{2}+r^{2}\right)^{-n / 2}(1+) e^{-t}+x\right)^{2}+r^{2}\right)^{-n / 2} d t d x d x_{r} . \tag{3.37}
\end{align*}
$$

As above, we can restrict to the case $x>1$. We get that this equals

$$
\begin{align*}
& \iiint_{0}^{\infty} y^{(n / 2)-1} r^{n-2}\left(1+x^{2}+r^{2}\right)^{-n / 2}\left(1+(x+y)^{2}+r^{2}\right)^{-n / 2} d y d x d r \\
& \quad=\iint_{0}^{\infty} \int_{x}^{\infty}(y-x)^{(n / 2)-1} r^{n-2}\left(1+x^{2}+r^{2}\right)^{-n / 2}\left(1+y^{2}+r^{2}\right)^{-n / 2} d y d x d r \tag{3.38}
\end{align*}
$$

which equals

$$
\begin{equation*}
\iint_{0}^{\infty} \int_{1}^{\infty}(v-1)^{n / 2-1} x^{n / 2} r^{n-2}\left(1+x^{2}+r^{2}\right)^{-n / 2}\left(1+v^{2} x^{2}+r^{2}\right)^{-n / 2} d y d x d r \tag{3.39}
\end{equation*}
$$

As above, it suffices to restrict the integration to the domain $x>1$. So one considers

$$
\begin{equation*}
\int_{0}^{\infty} \iint_{1}^{\infty}(v-1)^{(n / 2)-1} x^{n / 2} r^{n-2}\left(1+x^{2}+r^{2}\right)^{-n / 2}\left(1+v^{2} x^{2}+r^{2}\right)^{-n / 2} d y d x d r \tag{3.40}
\end{equation*}
$$

Choose $0<\varepsilon<1 / 2$ and write

$$
\begin{align*}
\left(1+x^{2}+r^{2}\right)^{-n / 2} & =\left(1+x^{2}+r^{2}\right)^{\varepsilon-n / 2}\left(1+x^{2}+r^{2}\right)^{-\varepsilon} \\
& \leq\left(1+r^{2}\right)^{\varepsilon-n / 2}\left(1+x^{2}\right)^{-\varepsilon} \tag{3.41}
\end{align*}
$$

So our integral is less than or equal to

$$
\begin{equation*}
\int_{0}^{\infty} r^{n-2}\left(1+r^{2}\right)^{\varepsilon-n / 2} d r \tag{3.42}
\end{equation*}
$$

times

$$
\begin{align*}
& \iint_{1}^{\infty}(v-1)^{(n / 2)-1} x^{n / 2}\left(1+x^{2}\right)^{-\varepsilon}\left(1+v^{2} x^{2}\right)^{-n / 2} d v d x \\
& \quad \leq C \iint_{1}^{\infty}(v-1)^{(n / 2)-1} x^{n / 2}\left(1+x^{2}\right)^{-\varepsilon}\left(1+v^{2}\right)^{-n / 2}\left(1+x^{2}\right)^{-n / 2} d v d x<\infty . \tag{3.43}
\end{align*}
$$

## 4. An explicit formula

Let $d \geq 2$ and $G=\mathrm{SO}(d, 1)^{0}$ if $d>2$. For $d=2$, let $G$ be the double cover $\widetilde{\mathrm{SO}}(2,1)^{o} \cong$ $\operatorname{PGL}_{2}(\mathbb{R})$. Then $K=\mathrm{SO}(d), M=\mathrm{SO}(d-1)$ for $d>2$. For $d=2$, we have $K \cong O(2)$ and $M \cong \mathbb{Z} / 2 \mathbb{Z}$ and in all cases, we have $M \backslash K \cong S^{d-1}$, the $(d-1)$-dimensional sphere. For each $\lambda \in \mathfrak{a}^{*}$, let $e_{\lambda}$ be the class-one vector in the associated principal series representation $\pi_{\lambda}$ given by $e_{\lambda}(a n k)=a^{\lambda+\rho}$. Let $\lambda, \mu, \nu \in \mathfrak{a}^{*}$ be imaginary. Let $\mathscr{T}_{\text {st }}$ be the invariant distribution on $L_{(\lambda, \mu, \nu)}$ considered in the last section. We are interested in the growth of $\mathscr{T}_{\text {st }}\left(e_{\lambda}, e_{\mu}, e_{\nu}\right)$ as a function in $\lambda$. First note that the Killing form induces a norm $|\cdot|$ on $\mathfrak{a}^{*}$.

We write $\mathscr{T}_{\text {st }}(\lambda, \mu, \nu)$ for $\mathscr{T}_{\text {st }}\left(e_{\lambda}, e_{\mu}, e_{\nu}\right)$ and identifying $\mathfrak{a}_{\mathbb{R}}$ to $\mathbb{R}$ via $\lambda \mapsto \lambda\left(H_{0}\right)$, we consider $\mathscr{T}_{\text {st }}$ as a function on $(i \mathbb{R})^{3}$.

In this section, we will prove the following theorem.
Theorem 4.1. For $\lambda, \mu$, $\nu$ imaginary, $\mathscr{T}_{\text {st }}(\lambda, \mu, \nu)$ equals a positive constant times

$$
\begin{equation*}
\frac{\Gamma((2(\lambda+\mu-\nu)+n) / 4) \Gamma((2(\lambda-\mu+\nu)+n) / 4) \Gamma((2(-\lambda+\mu+\nu)+n) / 4) \Gamma((2(\lambda+\mu+\nu)+n) / 4)}{\Gamma((2 \lambda+n) / 2) \Gamma((2 \mu+n) / 2) \Gamma((2 \nu+n) / 2)} \tag{4.1}
\end{equation*}
$$

where $n=d-1$, so in particular, for fixed imaginary $\mu$ and $\nu$. Then, as $|\lambda|$ tends to infinity, while $\lambda$ is imaginary, one has the asymptotic

$$
\begin{equation*}
\left|\mathscr{T}_{\text {st }}(\lambda, \mu, \nu)\right|=c \exp \left(-\frac{\pi}{2}|\lambda|\right)|\lambda|^{(d / 2)-2}\left(1+O\left(\frac{1}{|\lambda|}\right)\right) \tag{4.2}
\end{equation*}
$$

for some constant c $>0$.

Proof. The asymptotic formula follows from the explicit expression and the well-known asymptotical formula,

$$
\begin{equation*}
|\Gamma(\sigma+i t)|=\sqrt{2 \pi} \exp \left(-\frac{\pi}{2}|t|\right)|t|^{\sigma-(1 / 2)}\left(1+O\left(\frac{1}{|t|}\right)\right) \tag{4.3}
\end{equation*}
$$

as $|t| \rightarrow \infty$, where the real part $\sigma$ is fixed.
Now for the proof of the first assertion. The permutation group $S_{3}$ in three letters acts on $(i \mathbb{R})^{3}$ by permuting the coordinates. We claim that $\mathscr{T}_{\text {st }}$ is invariant under that action. So let $\sigma \in S_{3}$ and let $f=f_{1} \otimes f_{2} \otimes f_{3}$, then

$$
\begin{equation*}
\mathscr{T}_{\mathrm{st}}(\sigma(f))=\int_{G} f\left(\sigma^{-1}\left(x_{0} \cdot y\right)\right) d y=\int_{G} f\left(\sigma^{-1}\left(x_{0}\right) \cdot y\right) d y . \tag{4.4}
\end{equation*}
$$

Since the open orbit is unique, there is $y_{\sigma} \in G$ with $\sigma^{-1}\left(x_{0}\right)=x_{0} \cdot y_{\sigma}$, and so

$$
\begin{equation*}
\mathscr{T}_{\text {st }}(\sigma(f))=\int_{G} f\left(x_{0} \cdot y_{\sigma} y\right) d y=\int_{G} f\left(x_{0} \cdot y\right) d y=\mathscr{T}_{\text {st }}(f) . \tag{4.5}
\end{equation*}
$$

So in particular $\mathscr{T}_{\text {st }}(\lambda, \mu, \nu)$ is invariant under permutations of $(\lambda, \mu, \nu)$.
We have

$$
\begin{equation*}
\mathscr{T}_{\text {st }}\left(e_{\lambda}, e_{\mu}, e_{\nu}\right)=\int_{G} \underline{a}\left(w_{0} n_{0} y\right)^{\lambda+\rho} \underline{a}\left(w_{0} y\right)^{\mu+\rho} \underline{a}(y)^{\nu+\rho} d y . \tag{4.6}
\end{equation*}
$$

Replace the integral over $G$ by an integral over ANK using the Iwasawa decomposition. Since $\underline{a}(x k)=\underline{a}(x)$ for $x \in G$ and $k \in K$, the $K$-factor is irrelevant and we have to compute

$$
\begin{equation*}
\int_{A N} \underline{a}\left(w_{0} n_{0} a n\right)^{\lambda+\rho} \underline{a}\left(w_{0} a n\right)^{\mu+\rho} a^{\nu+\rho} d a d n . \tag{4.7}
\end{equation*}
$$

Note that $w_{0} a=a^{-1} w_{0}$, and so we have $\underline{a}\left(w_{0} a n\right)^{\rho}=a^{-\rho} \underline{a}\left(w_{0} n\right)$ as well as $\underline{a}\left(w_{0} n_{0} a n\right)=$ $a^{-\rho} \underline{a}\left(w_{0} n_{0}^{a} n\right)$, where $n_{0}^{a}=a^{-1} n_{0} a$. So the integral equals

$$
\begin{equation*}
\int_{A N} a^{\nu-\lambda-\mu-\rho} \underline{a}\left(w_{0} n_{0}^{a} n\right)^{\lambda+\rho} \underline{a}\left(w_{0} n\right)^{\mu+\rho} d a d n \tag{4.8}
\end{equation*}
$$

Using the explicit approach of the last section, we see that $\mathscr{T}_{\text {st }}(\lambda, \mu, \nu)$ equals

$$
\begin{equation*}
0 \int_{\mathbb{R}} \int_{\mathbb{R}^{d-1}} e^{t(v-\lambda-\mu-(d-1) / 2)}\left(1+\left|e^{-t} v_{1}+x\right|^{2}\right)^{-\lambda-(d-1) / 2}\left(1+|x|^{2}\right)^{-\mu-(d-1) / 2} d x d t \tag{4.9}
\end{equation*}
$$

where $v_{1}=(1,0, \ldots, 0)$.
Set $n=d-1=1,2,3, \ldots$ and use polar coordinates to compute for $n>1$ that

$$
\begin{equation*}
\int_{\mathbb{R}} \int_{\mathbb{R}^{n}} e^{t(v-\lambda-\mu-(n / 2))}\left(1+\left|e^{-t} v_{1}+x\right|^{2}\right)^{-\lambda-(n / 2)}\left(1+|x|^{2}\right)^{-\mu-(n / 2)} d x d t \tag{4.10}
\end{equation*}
$$

equals $(n-1) \operatorname{vol}\left(B_{n-1}\right)$ times

$$
\begin{equation*}
\iint_{\mathbb{R}} \int_{0}^{\infty} r^{n-2} e^{t(\nu-\lambda-\mu-(n / 2))}\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-\lambda-(n / 2)}\left(1+x^{2}+r^{2}\right)^{-\mu-(n / 2)} d r d x d t \tag{4.11}
\end{equation*}
$$

Define $I_{n}(a, b, c)$ to be

$$
\begin{equation*}
\iint_{\mathbb{R}} \int_{0}^{\infty} r^{n-2} e^{t(c-a-b)}\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-a}\left(1+x^{2}+r^{2}\right)^{-b} d r d x d t \tag{4.12}
\end{equation*}
$$

and $I_{1}(a, b, c)$ to be equal to

$$
\begin{equation*}
\iint_{\mathbb{R}} e^{t(c-a-b)}\left(1+\left(e^{-t}+x\right)^{2}\right)^{-a}\left(1+x^{2}\right)^{-b} d r d x d t \tag{4.13}
\end{equation*}
$$

Then, for all $n \in \mathbb{N}$,

$$
\begin{equation*}
\mathscr{T}_{\text {st }}(\lambda, \mu, \nu)=c_{n} I_{n}\left(\lambda+\frac{n}{2}, \mu+\frac{n}{2}, \nu+\frac{n}{2}\right), \tag{4.14}
\end{equation*}
$$

where $c_{1}=1$ and $c_{n}=(n-1) \operatorname{vol}\left(B_{n-1}\right)$ for $n>1$. So in particular, $I_{n}$ is invariant under permutations of the arguments.

Note that

$$
\begin{equation*}
\frac{\partial}{\partial r}\left[\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-a}\left(1+x^{2}+r^{2}\right)^{-b}\right] \tag{4.15}
\end{equation*}
$$

equals

$$
\begin{equation*}
-2 a r\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-a-1}\left(1+x^{2}+r^{2}\right)^{-b}-2 b r\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-a}\left(1+x^{2}+r^{2}\right)^{-b-1} \tag{4.16}
\end{equation*}
$$

and $r^{n-2}=\partial / \partial r\left[r^{n-1} / n-1\right]$. So, integrating by parts, for $n>1$ we compute

$$
\begin{equation*}
I_{n}(a, b, c)=\frac{2 a}{n-1} I_{n+2}(a+1, b, c+1)+\frac{2 b}{n-1} I_{n+2}(a, b+1, c+1) . \tag{4.17}
\end{equation*}
$$

To get a similar result for $I_{1}$, note that

$$
\begin{equation*}
\left(1+\left(e^{-t}+x\right)^{2}\right)^{-a}\left(1+x^{2}\right)^{-b} \tag{4.18}
\end{equation*}
$$

equals

$$
\begin{align*}
-(1 & \left.+\left(e^{-t}+x\right)^{2}+r^{2}\right)\left.^{-a}\left(1+x^{2}+r^{2}\right)^{-b}\right|_{r=0} ^{r=\infty} \\
& =-\int_{0}^{\infty} \frac{\partial}{\partial r}\left[\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-a}\left(1+x^{2}+r^{2}\right)^{-b}\right] d r . \tag{4.19}
\end{align*}
$$

This implies that

$$
\begin{equation*}
I_{1}(a, b, c)=2 a I_{3}(a+1, b, c+1)+2 b I_{3}(a, b+1, c+1) . \tag{4.20}
\end{equation*}
$$

Let $I_{n}^{\prime}(a, b, c)$ be the same as $I_{n}(a, b, c)$ except that there is a factor $x$ in the integrand, that is,

$$
\begin{equation*}
\iint_{\mathbb{R}} \int_{0}^{\infty} r^{n-2} e^{t(c-a-b)}\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-a}\left(1+x^{2}+r^{2}\right)^{-b} x d r d x d t . \tag{4.21}
\end{equation*}
$$

Replacing $x$ with $-x$ first and then with $x+e^{-t} y$ yields

$$
\begin{equation*}
I_{n}^{\prime}(a, b, c)=-I_{n}^{\prime}(b, a, c)-I_{n}(b, a, c-1) . \tag{4.22}
\end{equation*}
$$

The last equation also holds for $n=1$. Integration by parts gives

$$
\begin{align*}
& \int_{\mathbb{R}} e^{t(c-a-b)}\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-a} d t \\
& \quad=\frac{2 a}{a+b-c} \int_{\mathbb{R}} e^{t(c-a-b)}\left(1+\left(e^{-t}+x\right)^{2}+r^{2}\right)^{-a-1}\left(e^{-2 t}+x e^{-t} v\right) d t \tag{4.23}
\end{align*}
$$

So,

$$
\begin{equation*}
I_{n}(a, b, c)=\frac{2 a}{a+b-c}\left(I_{n}(a+1, b, c-1)+I_{n}^{\prime}(a+1, b, c)\right) . \tag{4.24}
\end{equation*}
$$

Using (4.22) yields

$$
\begin{equation*}
I_{n}(a, b, c)=\frac{2 a}{c-a-b} I_{n}^{\prime}(b, a+1, c) \tag{4.25}
\end{equation*}
$$

or

$$
\begin{equation*}
I_{n}^{\prime}(x, y, z)=\frac{z-x-y+1}{2 y-2} I_{n}(y-1, x, z) . \tag{4.26}
\end{equation*}
$$

Plugging this into (4.22), one gets

$$
\begin{equation*}
\frac{c-a-b+1}{2 b-2} I_{n}(b-1, a, c)=\frac{a+b-c-1}{2 a-2} I_{n}(a-1, b, c)-I_{n}(b, a, c-1) . \tag{4.27}
\end{equation*}
$$

Replacing ( $a, b, c$ ) with $(a+1, b+1, c+1)$ gives

$$
\begin{equation*}
\frac{c-a-b}{2 b} I_{n}(b, a+1, c+1)=\frac{a+b-c}{2 a} I_{n}(a, b+1, c+1)-I_{n}(b+1, a+1, c), \tag{4.28}
\end{equation*}
$$

or, using the invariance of $I_{n}$,

$$
\begin{equation*}
\frac{c-b-a}{2 a b}\left(a I_{n}(a+1, b, c+1)+b I_{n}(a, b+1, c+1)\right)=-I_{n}(a+1, b+1, c), \tag{4.29}
\end{equation*}
$$

and hence

$$
\begin{equation*}
a I_{n+2}(a+1, b, c+1)+b I_{n+2}(a, b+1, c+1)=\frac{2 a b}{a+b-c} I_{n+2}(a+1, b+1, c) . \tag{4.30}
\end{equation*}
$$

Using (4.17) and (4.20), one gets for $n>1$ that

$$
\begin{gather*}
I_{n}(a, b, c)=\frac{4 a b}{a+b-c} \frac{1}{n-1} I_{n+2}(a+1, b+1, c),  \tag{4.31}\\
I_{1}(a, b, c)=\frac{4 a b}{a+b-c} I_{3}(a+1, b+1, c) .
\end{gather*}
$$

Finally we arrive at

$$
\begin{equation*}
I_{n+2}(a+1, b+1, c+1)=\frac{a+b-c-1}{4 a b}(n-1) I_{n}(a, b, c+1) \tag{4.32}
\end{equation*}
$$

for $n>1$ and

$$
\begin{equation*}
I_{3}(a+1, b+1, c+1)=\frac{a+b-c-1}{4 a b} I_{1}(a, b, c+1) . \tag{4.33}
\end{equation*}
$$

These formulae amount to

$$
\begin{equation*}
\mathscr{T}_{\mathrm{st}}^{n+2}(\lambda, \mu, \nu)=d_{n} \frac{\lambda+\mu-\nu+(n / 2)-1}{(2 \lambda+n)(2 \mu+n)} \mathscr{T}_{\mathrm{st}}^{n}(\lambda, \mu, \nu+1), \tag{4.34}
\end{equation*}
$$

where $d_{n}=(n-1) c_{n+2} / c_{n}$ if $n>1$ and $d_{1}=c_{3}$. A calculation using the functional equation of the Gamma function, $\Gamma(z+1)=z \Gamma(z)$, shows that the right-hand side of the claim in the theorem satisfies the same equation as $n$ is replaced by $n+2$. So the claim of the theorem for $n$ implies the same claim for $n+2$. Note that [ 1 , formula (A.5)] implies the theorem for $n=1$, where one has to take into account that in [1], a different normalization is used. To get our formula from theirs, one has to replace $\lambda_{j}$ by $-2 \lambda_{j}$ in [1]. To finish the proof of the theorem, it therefore remains to show the claim for $n=2$.

To achieve this, we proceed in a fashion similar to [1]. First note that the group $G=$ $\mathrm{SL}_{2}(\mathbb{C})$ is a double cover of $\mathrm{SO}(3,1)^{0}$, so we might as well use this group for the computation. For $\lambda \in \mathbb{C}$, let $V_{\lambda}$ denote the space of all smooth functions $f$ on $\mathbb{C}^{2}$ with $f(a z, a w)=$ $|a|^{-2(\lambda+1)} f(z, w)$ for every $a \in \mathbb{C}^{\times}$. Then $G=\mathrm{SL}_{2}(\mathbb{C})$ acts on the space $V_{\lambda}$ via $\pi_{\lambda}(g) f(z, w)=$ $f((z, w) g)$. This is the principal series representation with parameter $\lambda$. For $\lambda=1$, there is a $G$-invariant continuous linear functional $\mathfrak{L}: V_{1} \rightarrow \mathbb{C}$, which is unique up to scalars and is given by

$$
\begin{equation*}
\mathfrak{L}(f)=\int_{S^{3}} f(x) d x \tag{4.35}
\end{equation*}
$$

where the integral runs over the standard sphere $S^{3} \subset \mathbb{C}^{2}$ with the volume element induced by the standard metric on $\mathbb{C}^{2} \cong \mathbb{R}^{4}$. Since $S^{3}$ equals $(0,1) K$, where $K=\operatorname{SU}(2)$ is the maximal compact subgroup of $G$, the theory of induced representations shows that this functional is indeed invariant under $G$ and is unique with that property. If $\lambda$ is imaginary, then the representation $\pi_{\lambda}$ is preunitary, the inner product being given by $\langle f, g\rangle=\mathfrak{L}(f \bar{g})$. To compute $\mathscr{T}_{\text {st }}^{2}\left(\lambda_{1}, \lambda_{2}, \lambda_{3}\right)$, we will describe this functional on the space $V_{\lambda_{1}} \otimes V_{\lambda_{2}} \otimes V_{\lambda_{3}}$. Let $\omega: \mathbb{C}^{2} \rightarrow \mathbb{C}$ be given by $\omega(a, b)=a_{1} b_{2}-a_{2} b_{1}$ the bilinear $G$-invariant on $\mathbb{C}^{2}$. Let $K_{\lambda_{1}, \lambda_{2}, \lambda_{3}}\left(v_{1}, v_{2}, v_{3}\right)$ be the invariant kernel

$$
\begin{equation*}
\left|\omega\left(v_{1}, v_{2}\right)\right|^{\left(-\lambda_{1}+\lambda_{2}+\lambda_{3}-1\right)}\left|\omega\left(v_{1}, v_{3}\right)\right|^{\left(\lambda_{1}-\lambda_{2}+\lambda_{3}-1\right)}\left|\omega\left(v_{2}, v_{3}\right)\right|^{\left(\lambda_{1}+\lambda_{2}-\lambda_{3}-1\right)} . \tag{4.36}
\end{equation*}
$$

This kernel is invariant under the diagonal action of $G$ and is homogeneous of degree $2\left(\lambda_{j}-1\right)$ with respect to the variable $v_{j}$. From this point, the computations run like in [1, formulae (A.4) and (A.5)]. The theorem follows.

## Appendix

## Automorphic coefficients

Let $\Gamma \subset G$ be a uniform lattice and consider the right regular representation of $G$ on $L^{2}(\Gamma \backslash G)$, which decomposes as a direct sum,

$$
\begin{equation*}
L^{2}(\Gamma \backslash G)=\bigoplus_{\pi \in \hat{G}} L^{2}(\Gamma \backslash G)(\pi) \tag{A.1}
\end{equation*}
$$

where the isotypic component $L^{2}(\Gamma \backslash G)(\pi)$ is zero for $\pi$ outside a countable set of $\pi \in \widehat{G}$ and is always isomorphic to a sum of finitely many copies of $\pi$. By the Sobolev lemma, one has $L^{2}(\Gamma \backslash G)^{\infty} \subset C^{\infty}(\Gamma \backslash G)$. A function $\varphi$ is called pure if $\varphi \in L^{2}(\Gamma \backslash G)(\pi)$ for some $\pi \in \hat{G}$. For three pure smooth $\varphi_{1}, \varphi_{2}, \varphi_{3}$, the integral $\int_{\Gamma \backslash G} \varphi_{1}(x) \varphi_{2}(x) \varphi_{3}(x) d x$ exists. Moreover, fix $\pi_{1}, \pi_{2}, \pi_{3} \in \hat{G}$ and fix $G$-equivariant embeddings $\sigma_{j}: V_{\pi_{j}} \rightarrow L^{2}(\Gamma \backslash G)$, then

$$
\begin{equation*}
\mathscr{T}_{\text {aut }}\left(v_{1}, v_{2}, v_{3}\right)=\int_{\Gamma \backslash G} \sigma_{1}\left(v_{1}\right)(x) \sigma_{2}\left(v_{2}\right)(x) \sigma_{3}\left(v_{3}\right)(x) d x \tag{A.2}
\end{equation*}
$$

defines a $G$-invariant continuous trilinear form on $V_{\pi_{1}}^{\infty} \times V_{\pi_{2}}^{\infty} \times V_{\pi_{3}}^{\infty}$. Let $L^{2}(\Gamma \backslash G)^{K}$ be the space of all $K$-invariant vectors in $L^{2}(\Gamma \backslash G)$. Then there is an orthonormal basis $\left(\varphi_{i}\right)_{i \in \mathbb{N}}$ of pure vectors in $L^{2}(\Gamma \backslash G)^{K}$. For each $i$, let $L_{i}$ be the $G$-stable closed subspace of $L^{2}(\Gamma \backslash G)$ generated by $\varphi_{i}$. Then the spaces $L_{i}$ are mutually orthogonal.

For fixed pure normalized $\varphi, \varphi^{\prime} \in L^{2}(\Gamma \backslash G)^{K}$, we are interested in the growth of the sequence

$$
\begin{equation*}
c_{i} \stackrel{\text { def }}{=} \mathscr{T}_{\text {aut }}\left(\varphi_{i}, \varphi, \varphi^{\prime}\right) . \tag{A.3}
\end{equation*}
$$

Let $\left(\lambda_{i}, \mu, \nu\right)$ be the induction parameters of the representations given by $\left(\varphi_{i}, \varphi, \varphi^{\prime}\right)$. First note that the numbers $c_{i}$ decay exponentially in $\left|\lambda_{i}\right|$. So we define the renormalized sequence

$$
\begin{equation*}
b_{i} \stackrel{\text { def }}{=}\left|c_{i}\right|^{2} \exp \left(\pi\left|\lambda_{i}\right|\right) \tag{A.4}
\end{equation*}
$$

In [1], it is shown that for $d=2$, one has

$$
\begin{equation*}
\sum_{\left|\lambda_{i}\right| \leq T} b_{i} \leq C \log T \tag{A.5}
\end{equation*}
$$

and in [8, Theorem 7.6], for $d \geq 3$,

$$
\begin{equation*}
\sum_{\left|\lambda_{i}\right| \leq T} b_{i} \leq C T^{2(d-2)} \tag{A.6}
\end{equation*}
$$

We are going to reprove this result for $d \geq 4$.

Theorem A.1. Let $d \geq 4$. There is $C>0$ such that for every $T>0$,

$$
\begin{equation*}
\sum_{\left|\lambda_{i}\right| \leq T} b_{i} \leq C T^{2(d-2)} \tag{A.7}
\end{equation*}
$$

Proof. By the uniqueness of trilinear forms, there exists a constant $a_{i}$ such that

$$
\begin{equation*}
\mathscr{T}_{\text {aut }}\left(\varphi_{i}, \varphi, \varphi^{\prime}\right)=a_{i} \mathscr{T}_{\text {st }}\left(\varphi_{i}, \varphi, \varphi^{\prime}\right) \tag{A.8}
\end{equation*}
$$

This constant depends on the embeddings of $\pi_{\lambda}, \pi_{\nu}$, $\pi_{\mu}$ into $L^{2}(\Gamma \backslash G)$, but we can normalize these embeddings by insisting that the standard class-one vectors $e_{\lambda}, e_{\mu}, e_{\nu}$ are mapped to $\varphi_{i}, \varphi, \varphi^{\prime}$. Thus the number $a_{i}$ only depends on $\varphi_{i}, \varphi, \varphi^{\prime}$.

Fix two smooth preunitary $G$-representations $V, V^{\prime}$ and let $L, L^{\prime}$ be their Hilbert completions. Consider the unitary representation of $G \times G$ on $L \otimes L^{\prime}$ and let $E=\left(L \otimes L^{\prime}\right)^{\infty}$ be its smooth part. Denote by $\mathscr{H}(E)$ the real vector space of all Hermitian forms on $E$ and denote by $\mathscr{H}^{+}(E)$ the cone of nonnegative forms.

Let $W$ be a smooth preunitary admissible representation of $G$. A $G$-invariant functional $\mathscr{T}: W \otimes V \otimes V^{\prime} \rightarrow \mathbb{C}$ defines a $G$-equivariant linear map $l_{\mathcal{T}}: V \otimes V^{\prime} \rightarrow W^{*}$ which extends to a $G$-map $l_{\mathcal{F}}: E \rightarrow \bar{W}$, where $\bar{W}$ is the space $W$ with complex conjugate linear structure. The Hermitian form $H_{W}$ on $W$ induces a form $H_{\bar{W}}$ on $\bar{W}$. Define a Hermitian form $H_{\mathcal{T}}$ on $E$ by $H_{\mathcal{T}}=l_{\mathscr{F}}^{*} H_{\bar{W}}$, that is, for $u, v \in E$, one has $H_{\mathscr{T}}(u, v)=H_{\bar{W}}\left(l_{\mathcal{F}}(u), l_{\mathcal{T}}(v)\right)$. Finally, suppose that $W=V_{\lambda}, V, V^{\prime}$ are all principal series representations and let $\mathscr{T}_{\text {st }}$ be the standard trilinear form, then we write $H_{\lambda}^{\text {st }}$ for the Hermitian form induced on $E$. We also assume that $V_{\lambda}, V, V^{\prime}$ are cuspidal, that is, those fixed embeddings into $L^{2, \infty}$ are given. Then the cuspidal trilinear form $\mathscr{T}$ induces a Hermitian form $H_{\lambda}^{\text {aut }}$ on $E$.

Next consider the space $L^{2, \infty}(\Gamma \backslash G \times \Gamma \backslash G)$. Let $H_{\Delta}$ denote the Hermitian form on $L^{2, \infty}(\Gamma \backslash$ $G \times \Gamma \backslash G)$ given by restriction to the diagonal, that is,

$$
\begin{equation*}
H_{\Delta}(u, v)=\int_{\Gamma \backslash G} u(x, x) \overline{v(x, x)} d x . \tag{A.9}
\end{equation*}
$$

Let $U$ be a Hilbert space with Hermitian form $H$. For a closed subspace $L$ of $U$, let $P r_{L}$ denote the orthogonal projection from $U$ to $L$ and let $H_{L}$ be the Hermitian form on $U$ given by

$$
\begin{equation*}
H_{L}(u, v)=H\left(\operatorname{Pr}_{L}(u), \operatorname{Pr}_{L}(v)\right) . \tag{A.10}
\end{equation*}
$$

The map $L \mapsto H_{L}$ is additive and monotonic, that is, $H_{L+L^{\prime}}=H_{L}+H_{L^{\prime}}$ if $L$ and $L^{\prime}$ are orthogonal and $H_{L} \leq H_{L^{\prime}}$ if $L \subset L^{\prime}$.

Lemma A.2. On the space $E$, one has

$$
\begin{equation*}
\sum_{i}\left|a_{i}\right|^{2} H_{\lambda_{i}}^{\mathrm{st}} \leq H_{\Delta} . \tag{A.11}
\end{equation*}
$$

Proof. By the uniqueness, it follows that $H_{\lambda_{i}}^{\text {aut }}=\left|a_{i}\right|^{2} H_{\lambda_{i}}^{\text {st }}$ and since all the spaces $L_{i}$ are orthogonal, one also deduces that $\sum_{i} H_{\lambda_{i}}^{\text {aut }} \leq H_{\Delta}$.

The group $G \times G$ acts on $L^{2, \infty}(\Gamma \backslash G \times \Gamma \backslash G)$ and thus on the real vector space of Hermitian forms on $L^{2, \infty}(\Gamma \backslash G \times \Gamma \backslash G)$. By integration, this action extends to a representation of the convolution algebra $C_{c}^{\infty}(G \times G)$. Note that a nonnegative function $h \in C_{c}^{\infty}(G \times G)$ will preserve to cone of nonnegative Hermitian forms.

Lemma A.3. Let $h \in C_{c}^{\infty}(G \times G)$ be nonnegative. Then there exists $C>0$ such that

$$
\begin{equation*}
h \cdot H_{\Delta} \leq C H_{\mathrm{aut}}, \tag{A.12}
\end{equation*}
$$

where $H_{\text {aut }}$ is the $L^{2}$-Hermitian form on $L^{2, \infty}(\Gamma \backslash G \times \Gamma \backslash G)$.
Proof. The form $h \cdot H_{\Delta}$ is given by integration over a smooth measure on $\Gamma \backslash G \times \Gamma \backslash G$. Being smooth, this measure is bounded by a multiple of the invariant measure.

A positive functional on $\mathscr{H}(E)$ is an additive map $\rho: \mathscr{H}^{+}(E) \rightarrow \mathbb{R}_{\geq 0} \cup\{\infty\}$. For example, every vector $u \in E$ gives rise to a positive functional $\delta_{u}$ defined as

$$
\begin{equation*}
\delta_{u}(H)=H(u, u) . \tag{A.13}
\end{equation*}
$$

Every positive functional is monotonic and homogeneous, that is,

$$
\begin{align*}
& \rho(H) \leq \rho\left(H^{\prime}\right) \quad \text { if } H \leq H^{\prime}, \\
& \rho(t H)=t \rho(H) \quad \text { for } t>0 . \tag{A.14}
\end{align*}
$$

Let $\rho$ be a positive functional and let $h_{\rho}(\lambda)=\rho\left(H_{\lambda}^{\text {st }}\right)$. Then Lemma A. 2 implies that

$$
\begin{equation*}
\sum_{i} h_{\rho}\left(\lambda_{i}\right)\left|a_{i}\right|^{2} \leq \rho\left(H_{\Delta}\right) . \tag{A.15}
\end{equation*}
$$

Proposition A.4. There are constants $T_{0}, C>0$ such that for every $T \geq T_{0}$, there is a positive functional $\rho_{T}$ on $\mathscr{H}(E)$ with

$$
\begin{equation*}
\rho_{T}\left(H_{\Delta}\right) \leq C T^{\operatorname{dim}(M \backslash K)+1} \tag{A.16}
\end{equation*}
$$

and with $h_{T}=h_{\rho_{T}}$ one has $h_{T}(\lambda) \geq 1 / C$ for every $|\lambda| \leq 2 T$.
Proof. Let $D$ be a relatively compact open subset of $G$ with $g \in D_{1} \Rightarrow g^{-1} \in D$. Assume further that $D$ is invariant under right and left translations by elements of $K$. Let $h \in$ $C_{c}^{\infty}(G \times G)$ such that $h$ is constantly one on $D \times D$ and such that $h$ is invariant under left and right translates by elements of $K$. The space $E$ can be identified with $C^{\infty}(M \backslash K \times$ $M \backslash K)$. Fix a norm $\|\cdot\|$ on the Lie algebra of $A N$.

Note that the map

$$
\begin{gather*}
\psi: A \times N \times M \longrightarrow P \backslash G \times P \backslash G, \\
(a, n, m) \longmapsto\left(w_{0} a n m, w_{0} n_{0} a n m\right) \tag{A.17}
\end{gather*}
$$

is a diffeomorphism onto an open subset of $(P \backslash G)^{2} \cong(M \backslash K)^{2}$. Let $F=\psi(1 \times 1 \times M)=$ the $M$-orbit of the point $\left(w_{0}, w_{0} n_{0}\right)$. Let $C_{1}, T>0$ be large and set

$$
\begin{equation*}
S=\left\{\psi(a, n, m):\|\log a n\|<\frac{1}{C_{1} T}, m \in M\right\} . \tag{A.18}
\end{equation*}
$$

Let $u_{T}$ be a nonnegative smooth function on $(M \backslash K)^{2}$ with support in $S$ such that

$$
\begin{equation*}
\int u d x d y=1, \quad \int|u|^{2} d x d y \leq T^{\operatorname{dim}(M \backslash K)+1} \tag{A.19}
\end{equation*}
$$

It is possible to construct such a function since the codimension of $F$ equals $\operatorname{dim}(M \backslash K)+$ 1. With these data, define a positive functional $\rho_{T}$ on $\mathscr{H}(E)$ by

$$
\begin{equation*}
\rho_{T}(H)=\delta_{u_{T}}(h \cdot H) \tag{A.20}
\end{equation*}
$$

Then one has

$$
\begin{equation*}
\rho_{T}\left(H_{\Delta}\right)=\rho_{u}\left(h \cdot H_{\Delta}\right) \leq C^{\prime} \rho_{u}\left(H_{s t}\right) \leq C T^{\operatorname{dim}(M \backslash K)+1} . \tag{A.21}
\end{equation*}
$$

Recall the standard trilinear form $\mathscr{T}_{\text {st }}$ on $V_{\lambda} \otimes V \otimes V^{\prime}$ and the corresponding linear map $l_{\mathscr{T}_{\text {st }}}: V \otimes V^{\prime} \rightarrow \bar{V}_{\lambda}$. We identify the space $\bar{V}_{\lambda}$ with $C^{\infty}(M \backslash K)$ and let $z$ denote the point $M 1 \in M \backslash K$. The Dirac measure $\delta_{z}$ at $z$ is a continuous linear functional on $C^{\infty}(M \backslash K)$. We get an induced Hermitian form $H_{z}$ on $E$ defined by

$$
\begin{equation*}
H_{z}(u, v)=l_{\mathscr{J}_{s t}}(u)(z) \overline{\mathscr{F}_{s t}(v)(z)} \tag{A.22}
\end{equation*}
$$

Lemma A.5. On E, one has

$$
\begin{equation*}
H_{\lambda}^{\mathrm{st}}=\int_{K}(k, k) \cdot H_{z} d k \tag{A.23}
\end{equation*}
$$

Proof. This follows from the fact that the invariant Hermitian form on $\bar{V}_{\lambda}$ is equal to $\int_{K} \pi(k) \tilde{H}_{z} d k$, where $\tilde{H}_{z}$ is the Hermitian form on $\bar{V}_{\lambda}$ given by $\tilde{H}_{z}(u, v)=u(z) \overline{v(z)}$.

Since the test function $h$ was assumed to be $K \times K$-invariant, it follows that

$$
\begin{equation*}
h_{T}(\lambda)=\rho_{T}\left(H_{\lambda}^{\mathrm{st}}\right)=\delta_{u}\left(h \cdot H_{z}\right) \tag{A.24}
\end{equation*}
$$

So in order to establish a lower bound for $\rho_{T}\left(H_{\lambda}^{\text {st }}\right)$, it suffices to give a lower bound for $\delta_{u}\left(x \cdot H_{z}\right)=\left|\left\langle\pi(x) f_{z}, u\right\rangle\right|^{2}$ for $x \in D$, where $f_{z}$ is the linear functional on $E$ given by $f_{z}(u)=\delta_{z}\left(l_{\mathcal{T}_{\mathrm{st}}}(u)\right)$.

Lemma A.6. There is $C>0$ such that for $T_{0}$ large enough, there exists an open, nonempty subset $D_{0} \subset D$ such that for $T \geq T_{0},|\lambda| \leq T$, and $x \in D_{0},\left|\left\langle\pi(x) f_{z}, u_{T}\right\rangle\right| \geq 1 / C$.

Proof. Note first that with $x=a n k \in G$, we have

$$
\begin{equation*}
\left\langle\pi(x) f_{z}, u_{T}\right\rangle=\int_{(M \backslash K)^{2}} u_{T}\left(k_{2}, k_{3}\right) \phi\left(k, k_{2}^{a n k}, k_{3}^{a n k}\right) d k_{2} d k_{3} . \tag{A.25}
\end{equation*}
$$

We next derive an explicit formula for $\phi$ on the open orbit. Recall that

$$
\begin{align*}
\mathscr{T}_{\text {st }}(f) & \left.=\int_{G} f(y, w) y, w_{0} n_{0} y\right) d y \\
& =\int_{A N K} f\left(a n k, w_{0} a n k, w_{0} n_{0} a n k\right) d a d n d k \tag{A.26}
\end{align*}
$$

and that the last integrand equals

$$
\begin{equation*}
a^{\lambda+\rho} \underline{a}\left(w_{0} a n\right)^{\mu+\rho} \underline{a}\left(w_{0} n_{0} a n\right)^{v+\rho} f\left(k, \underline{k}\left(w_{0} a n\right) k, \underline{k}\left(w_{0} n_{0} a n\right) k\right) d a d n d k . \tag{A.27}
\end{equation*}
$$

We deduce that

$$
\begin{equation*}
\phi\left(k, \underline{k}\left(w_{0} a n\right) k, \underline{k}\left(w_{0} n_{0} a n\right) k\right)=a^{\lambda+\rho} \underline{a}\left(w_{0} a n\right)^{\mu+\rho} \underline{a}\left(w_{0} n_{0} a n\right)^{\nu+\rho} W(a n), \tag{A.28}
\end{equation*}
$$

where $W$ is the Radon-Nykodim derivative of $d a d n d k$ against the measure on $(M \backslash K)^{3}$. In particular, this function is positive and smooth and does not depend on $\lambda$. It follows that the functional $\pi(x) f_{z}$ is given by integration against the function

$$
\begin{equation*}
\pi(a n k) f_{z}\left(w_{0} a^{\prime} n^{\prime} m^{\prime}, w_{0} n_{0} a^{\prime} n^{\prime} m^{\prime}\right)=\phi\left(k, w_{0} a a^{\prime} n^{\prime \prime} m^{\prime} k, w_{0} n_{0} a a^{\prime} n^{\prime \prime} m^{\prime} k\right), \tag{A.29}
\end{equation*}
$$

where $n^{\prime \prime}=a^{-1} n^{\prime} a m^{\prime} n\left(m^{\prime}\right)^{-1}$. Note that $M k=M m^{\prime} k$ and so the above equals

$$
\begin{equation*}
\left(a a^{\prime}\right)^{\lambda+\rho} \underline{a}\left(w_{0} a a^{\prime} n^{\prime \prime}\right)^{\mu+\rho} \underline{a}\left(w_{0} n_{0} a a^{\prime} n^{\prime \prime}\right)^{\nu+\rho} W\left(a a^{\prime} n^{\prime \prime}\right) . \tag{A.30}
\end{equation*}
$$

Note in particular that the function $\pi(x) f_{z}$ is invariant under $M$.
The set $S$ is a tubular neighborhood of the compact set $F$. We have natural local coordinates with $F$-directions, which can be identified with open parts of $M$-orbits and $F^{\perp}$-directions, which can be viewed as $A N$-orbits. The above formula shows that the gradient of $\pi(x) f_{z}$ is zero along the $M$-orbits. Along the $A N$-orbits, this gradient is bounded by a constant times $T$, if we leave $\mu$ and $\nu$ fixed and restrict to $|\lambda| \leq T$. Let $C_{0}>0$ be large and let $D_{0} \subset D$ be the subset of elements $x \in D$ such that $\left|\pi(x) f_{z}\right|<C_{0}$. This set is open and nonempty if $C_{0}$ is large enough and it does not depend on $\lambda$.

Lemma A. 6 implies Proposition A. 4 and the latter implies Theorem A. 1 as follows.
Consider the inequality $\sum_{i}\left|a_{i}\right|^{2} \rho_{T}\left(H_{\lambda_{i}}^{\text {st }}\right) \leq \rho_{T}\left(H_{\Delta}\right)$. By Proposition A.4, the right-hand side is bounded by $C T^{2 \operatorname{dim}(M \backslash K)}$ and since $\operatorname{dim} M \backslash K=\operatorname{dim} P \backslash G=\operatorname{dim} N=d-1$, we get

$$
\begin{equation*}
\sum_{i}\left|a_{i}\right|^{2} \rho_{T}\left(H_{\lambda_{i}}^{\mathrm{st}}\right) \leq C T^{d} \tag{A.31}
\end{equation*}
$$

Restricting the sum to those $i$ with $\left|\lambda_{i}\right| \leq 2 T$ and using the second assertion of Proposition A.4, we get

$$
\begin{equation*}
\sum_{\left|\lambda_{i}\right| \leq T}\left|a_{i}\right|^{2} \leq C T^{d} . \tag{A.32}
\end{equation*}
$$

According to Theorem 4.1, we have $b_{i}\left|\lambda_{i}\right|^{4-d} \sim\left|a_{i}\right|^{2}$, and so

$$
\begin{equation*}
\sum_{\left|\lambda_{i}\right| \leq T} b_{i}\left|\lambda_{i}\right|^{4-d} \leq C T^{d} \tag{A.33}
\end{equation*}
$$

This implies the theorem.

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