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Lamara Bitsadze

**EXPLICIT SOLUTION OF THE FIRST BVP
OF THE ELASTIC MIXTURE FOR HALF-SPACE**

Abstract. We consider the first BVP of elastic mixture theory for a transversally-isotropic half-space. The solution of the first BVP for the transversally-isotropic half-space is given in [1]. The present paper is an attempt to use this result for the BVP of elastic mixture theory for a transversally-isotropic elastic body. Using the potential method and the theory of integral equations, the uniqueness theorem is proved for a half-space and the first BVP previously is solved effectively (in quadratures), which has not been solved.

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Տրված է երկրորդ կարգի գծափակ խառնուրդային առարկայի առաջին սահմանային խնդիրը առնչությամբ կիսադոմենում (կիսադոմենային համասեռ խառնուրդային առարկայի համար): Գտնվում է առաջին սահմանային խնդիրը լուծելու մեթոդը (կազմվում է ինտեգրալ հավասարումների միջոցով): Գտնվում է խնդիրը լուծելու մեթոդը (կազմվում է ինտեգրալ հավասարումների միջոցով): Գտնվում է խնդիրը լուծելու մեթոդը (կազմվում է ինտեգրալ հավասարումների միջոցով):

The first BVP and the uniqueness theorem for a half-space. Let the plane ox_1x_2 be the boundary of a half-space $x_3 > 0$. Let the upper half-space be denoted by D and the boundary of D by S . Let the axis ox_3 be directed vertically upwards and the normal be $n(0, 0, 1)$.

A basic homogeneous equation of statics of transversally-isotropic elastic mixture theory can be written in the form [2]

$$C(\partial x)U = \begin{pmatrix} C^{(1)}(\partial x) & C^{(3)}(\partial x) \\ C^{(3)}(\partial x) & C^{(2)}(\partial x) \end{pmatrix} U = 0, \quad (1)$$

where the components of the matrix $C^{(j)}(\partial x) = \|C_{pq}^{(j)}(\partial x)\|_{3 \times 3}$ are given in the form

$$\begin{aligned} C_{pq}^{(j)} &= C_{qp}^{(j)}, \quad j = 1, 2, 3; \quad p, q = 1, 2, 3, \\ C_{11}^{(j)}(\partial x) &= c_{11}^{(j)} \frac{\partial^2}{\partial x_1^2} + c_{66}^{(j)} \frac{\partial^2}{\partial x_2^2} + c_{44}^{(j)} \frac{\partial^2}{\partial x_3^2}, \\ C_{12}^{(j)}(\partial x) &= (c_{11}^{(j)} - c_{66}^{(j)}) \frac{\partial^2}{\partial x_1 \partial x_2}, \\ C_{k3}^{(j)}(\partial x) &= (c_{13}^{(j)} + c_{44}^{(j)}) \frac{\partial^2}{\partial x_k \partial x_3}, \quad k = 1, 2, \\ C_{22}^{(j)}(\partial x) &= c_{66}^{(j)} \frac{\partial^2}{\partial x_1^2} + c_{11}^{(j)} \frac{\partial^2}{\partial x_2^2} + c_{44}^{(j)} \frac{\partial^2}{\partial x_3^2}, \\ C_{33}^{(j)}(\partial x) &= c_{44}^{(j)} \left(\frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} \right) + c_{33}^{(j)} \frac{\partial^2}{\partial x_3^2}, \end{aligned}$$

$c_{pq}^{(k)}$ are the constants characterizing physical properties of the mixture and satisfying certain inequalities obtained due to positive definiteness of the potential energy. $U = U^T(x) = (u', u'')$ is a six-dimensional displacement vector-function, $u'(x) = (u'_1, u'_2, u'_3)$ and $u''(x) = (u''_1, u''_2, u''_3)$ are partial displacement vectors. Throughout this paper “ T ” denotes transposition.

Definition. A vector-function $U(x)$ defined in the domain D is called regular if it has integrable continuous second derivatives in D and $U(x)$ itself and its first derivatives are continuously extendable at every point of the boundary of D , i.e. $U(x) \in C^2(D) \cap C^1(D)$ and satisfies the following conditions at infinity

$$U(x) = O(|x|^{-1}), \quad \frac{\partial U}{\partial x_k} = O(|x|^{-2}), \quad |x|^2 = x_1^2 + x_2^2 + x_3^2, \quad k = 1, 2, 3.$$

For the equation (1) we pose the following BVP. Find a regular function $U(x)$ satisfying the equation (1) in D if on the boundary S the displacement vector U is given in the form

$$U^+ = f(z), \quad z \in S. \quad (2)$$

where $(\cdot)^+$ denotes the limiting value from D and f is a given vector.

$$\begin{aligned} |f_k| < AR, \quad R = \sqrt{z_1^2 + z_2^2} \leq 1, \quad |f_k| < AR^{-\alpha}, \\ \alpha > 0, \quad R > 1, \quad k = 1, \dots, 6, \quad A = \text{const} > 0. \end{aligned} \quad (3)$$

The Uniqueness Theorem. Let us prove that the first homogeneous BVP has only a trivial solution. Note that if U is a regular solution of the equation (1) and satisfies the following conditions at infinity

$$U(x) = O(|x|^{-\alpha}), \quad P(\partial x, n)U = O(|x|^{-1-\alpha}), \quad \alpha > 0,$$

then we have the formula

$$\begin{aligned} U(x) &= \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} [(P(\partial y, n)\Gamma)^* u^+ - \Gamma(y-z)(P(\partial y, n)u)^+] dy_1 dy_2, \quad x \in D, \end{aligned} \quad (4)$$

where $P(\partial y, n)U$ is the generalized stress vector

$$\begin{aligned} (P(\partial y, n)U)_k &= c_{44}^{(1)} \frac{\partial u'_k}{\partial x_3} + c_{44}^{(3)} \frac{\partial u''_k}{\partial x_3} + \delta^{(1)} \frac{\partial u'_3}{\partial x_k} + \delta^{(3)} \frac{\partial u''_3}{\partial x_k}, \quad k = 1, 2, \\ (P(\partial y, n)U)_3 &= \beta^{(1)} \left(\frac{\partial u'_1}{\partial x_1} + \frac{\partial u'_2}{\partial x_2} \right) + \beta^{(3)} \left(\frac{\partial u''_1}{\partial x_1} + \frac{\partial u''_2}{\partial x_2} \right) + \\ &\quad + c_{33}^{(1)} \frac{\partial u'_3}{\partial x_3} + c_{33}^{(3)} \frac{\partial u''_3}{\partial x_3}, \\ (P(\partial y, n)U)_k &= c_{44}^{(3)} \frac{\partial u'_{k-3}}{\partial x_3} + c_{44}^{(2)} \frac{\partial u''_{k-3}}{\partial x_3} + \\ &\quad + \delta^{(4)} \frac{\partial u'_3}{\partial x_{k-3}} + \delta^{(2)} \frac{\partial u''_3}{\partial x_{k-3}}, \quad k = 4, 5, \\ (P(\partial y, n)U)_6 &= \beta^{(4)} \left(\frac{\partial u'_1}{\partial x_1} + \frac{\partial u'_2}{\partial x_2} \right) + \beta^{(2)} \left(\frac{\partial u''_1}{\partial x_1} + \frac{\partial u''_2}{\partial x_2} \right) + \\ &\quad + c_{33}^{(3)} \frac{\partial u'_3}{\partial x_3} + c_{33}^{(2)} \frac{\partial u''_3}{\partial x_3}, \\ \beta^{(j)} + \delta^{(j)} &= \alpha_{13}^{(j)}, \quad j = 1, 2, 3, \quad \beta^{(4)} + \delta^{(4)} = \alpha_{13}^{(3)}, \\ c_{13}^{(j)} + c_{44}^{(j)} &= \alpha_{13}^{(j)}. \end{aligned} \quad (5)$$

$\Gamma(y-x)$ is the symmetric matrix of the fundamental solution of the equation (1)

$$\Gamma(x-y) = \begin{pmatrix} \Gamma^{(1)} & \Gamma^{(3)} \\ \Gamma^{(3)T} & \Gamma^{(2)} \end{pmatrix}, \quad (6)$$

where

$$\Gamma^{(j)}(x-y) = \sum_{k=1}^6 \|\Gamma_{pq}^{j(k)}\|_{3 \times 3}, \quad j = 1, 2, 3, \quad \Gamma_{pq}^{j(k)} = \Gamma_{qp}^{j(k)},$$

$$\begin{aligned}
\Gamma_{pq}^{1(k)} &= \delta_{pq} \frac{A_{11}^{(k)}}{r_k} + A_{12}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_q}, \quad p = 1, 2; \quad q = 1, 2; \\
\delta_{pq} &= 1, \quad p = q, \quad \delta_{pq} = 0, \quad p \neq q, \\
\Gamma_{p3}^{1(k)} &= A_{13}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3}, \quad \Gamma_{33}^{1(k)} = \frac{A_{33}^{(k)}}{r_k}, \quad \Gamma_{pq}^{3(k)} = \delta_{pq} \frac{A_{14}^{(k)}}{r_k} + A_{42}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_q}, \\
\Gamma_{p3}^{3(k)} &= A_{16}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3}, \quad p = 1, 2, \quad \Gamma_{33}^{3(k)} = \frac{A_{36}^{(k)}}{r_k}, \quad \Gamma_{3p}^{3(k)} = A_{34}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3}, \\
\Gamma_{pq}^{2(k)} &= \delta_{pq} \frac{A_{44}^{(k)}}{r_k} + A_{45}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_q}, \\
\Gamma_{p3}^{2(k)} &= A_{46}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3}, \quad p = 1, 2, \quad \Gamma_{33}^{2(k)} = \frac{A_{66}^{(k)}}{r_k}.
\end{aligned}$$

The coefficients $A_{pq}^{(k)}$ are defined as follows

$$\begin{aligned}
A_{11}^{(k)} &= (-1)^k (c_{44}^{(2)} - c_{66}^{(2)} a_k) r_0', \quad A_{14}^{(k)} = -(-1)^k (c_{44}^{(3)} - c_{66}^{(3)} a_k) r_0', \\
A_{12}^{(k)} &= \frac{A_{11}^{(k)}}{a_k}, \quad A_{24}^{(k)} = \frac{A_{14}^{(k)}}{a_k}, \quad A_{45}^{(k)} = \frac{A_{44}^{(k)}}{a_k}, \\
A_{44}^{(k)} &= (-1)^k (c_{44}^{(1)} - c_{66}^{(1)} a_k) r_0', \quad k = 1, 2, \quad r_0' = [r_0(a_1 - a_2)]^{-1}, \\
A_{12}^{(k)} &= \frac{\delta_k}{a_k} [-q_3 c_{44}^{(2)} + a_k t_{12} - a_k^2 t_{11} + c_{11}^{(2)} q_4 a_k^3], \\
A_{42}^{(k)} &= \frac{\delta_k}{a_k} [q_3 c_{44}^{(3)} + a_k t_{13} - a_k^2 t_{22} - c_{11}^{(3)} q_4 a_k^3], \\
A_{45}^{(k)} &= \frac{\delta_k}{a_k} [-q_3 c_{44}^{(1)} + a_k t_{23} - a_k^2 t_{33} + c_{11}^{(1)} q_4 a_k^3], \\
A_{33}^{((k))} &= \delta_k [q_4 c_{33}^{(2)} - a_k t_{42} + a_k^2 t_{44} - c_{44}^{(2)} q_1 a_k^3], \\
A_{36}^{(k)} &= \delta_k [-q_4 c_{33}^{(3)} - a_k t_{62} + a_k^2 t_{66} + c_{44}^{(3)} q_1 a_k^3], \\
A_{66}^{(k)} &= \delta_k [q_4 c_{33}^{(1)} - a_k t_{52} + a_k^2 t_{55} - c_{44}^{(1)} q_1 a_k^3], \\
A_{13}^{(k)} &= \delta_k [v_{13} - v_{11} a_k + v_{12} a_k^2], \quad A_{16}^{(k)} = \delta_k [w_{13} - w_{12} a_k + w_{11} a_k^2], \\
A_{34}^{(k)} &= \delta_k [v_{23} - v_{21} a_k + v_{22} a_k^2], \quad A_{46}^{(k)} = \delta_k [w_{34} - w_{14} a_k + w_{24} a_k^2], \\
\delta_k &= d_k (a_1 - a_k) (a_2 - a_k) b_0^{-1}, \quad k = 3, \dots, 6,
\end{aligned} \tag{7}$$

where a_k are the positive roots of the characteristic equations

$$\begin{aligned}
(r_0 a^2 - c_0 a + q_4)(b_0 a^4 - b_1 a^3 + b_2 a^2 - b_3 a + b_4) &= 0, \\
r_0 &= c_{66}^{(1)} c_{66}^{(2)} - c_{66}^{(3)2}, \quad c_0 = c_{66}^{(1)} c_{44}^{(2)} + c_{44}^{(1)} c_{66}^{(2)} - 2c_{66}^{(3)} c_{44}^{(3)}.
\end{aligned}$$

The coefficients d_k , b_k , v_{ij} , w_{ij} , t_{ij} are given in [3]. The singular matrix $[P(\partial y, n)\Gamma]^*$ = $\sum_{k=1}^6 (M_{pq}^{(k)})_{6 \times 6}$, which is obtained from $P(\partial x, n)\Gamma(x - y)$ by

transposition of the columns and rows and the variables x and y , has the form

$$[P(\partial x)\Gamma(x-y)]^* = \sum_{k=1}^6 \begin{pmatrix} M^{(1k)} & M^{(3k)} \\ M^{(4k)} & M^{(2k)} \end{pmatrix}, \quad (8)$$

where the elements of the matrix $M^{(jk)} = \|M_{pq}^{(jk)}\|_{3 \times 3}$, $j = 1, 2, 3, 4$, are written as

$$\begin{aligned} M_{pj}^{(1k)} &= \delta_{pj} R_{11}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k} + R_{12}^{(k)} \frac{\partial^3 \Phi_k}{\partial x_p \partial x_j \partial x_3}, \\ \delta_{pj} &= 1, \quad p = j, \quad \delta_{pj} = 0, \quad p \neq j, \quad p, j = 1, 2, \\ M_{p3}^{(1k)} &= R_{31}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \quad M_{3p}^{(1k)} = R_{13}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \\ M_{33}^{(1k)} &= R_{33}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k}, \quad M_{pj}^{(3k)} = \delta_{pj} R_{14}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k} + R_{24}^{(k)} \frac{\partial^3 \Phi_k}{\partial x_j \partial x_p \partial x_3}, \\ M_{p3}^{(3k)} &= R_{61}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \quad M_{3p}^{(3k)} = R_{43}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \quad M_{33}^{(3k)} = R_{63}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k}, \\ M_{pj}^{(4k)} &= \delta_{pj} R_{41}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k} + R_{42}^{(k)} \frac{\partial^3 \Phi_k}{\partial x_j \partial x_p \partial x_3}, \quad M_{p3}^{(4k)} = R_{34}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \\ M_{3p}^{(4k)} &= R_{16}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \quad M_{33}^{(4k)} = R_{36}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k}, \\ M_{pj}^{(2k)} &= \delta_{pj} \mu_{44}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k} + R_{44}^{(k)} \frac{\partial^3 \Phi_k}{\partial x_p \partial x_j \partial x_3}, \quad M_{p3}^{(2k)} = R_{64}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \\ M_{3p}^{(2k)} &= R_{46}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \quad M_{33}^{(2k)} = R_{66}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k}, \quad p = 1, 2. \end{aligned}$$

The coefficients $R_{pq}^{(k)}$ satisfy the following conditions

$$\begin{aligned} \sum_{k=1}^2 \frac{R_{11}^{(k)}}{a_k} &= \sum_{k=3}^6 \frac{R_{33}^{(k)}}{a_k} = \sum_{k=3}^6 \frac{R_{66}^{(k)}}{a_k} = \sum_{k=1}^2 \frac{\mu_{44}^{(k)}}{a_k} = 1, \\ \sum_{k=1}^2 \frac{R_{14}^{(k)}}{a_k} &= \sum_{k=1}^6 R_{12}^{(k)} = 0, \\ \sum_{k=1}^2 \frac{R_{41}^{(k)}}{a_k} &= \sum_{k=3}^6 \frac{R_{36}^{(k)}}{a_k} = \sum_{k=1}^6 R_{24}^{(k)} = \sum_{k=3}^6 \frac{R_{63}^{(k)}}{a_k} = \sum_{k=1}^2 R_{44}^{(k)} = \sum_{k=1}^6 R_{42}^{(k)} = 0 \end{aligned}$$

and, after elementary calculations the coefficients $R_{13}^{(k)}, \dots, R_{64}^{(k)}$ take the form

$$\begin{aligned}
R_{13}^{(k)} &= \delta_0^{(1)} A_{33}^{(k)} + \delta_0^{(3)} A_{36}^{(k)} + c_{44}^{(1)} A_{13}^{(k)} + c_{44}^{(3)} A_{43}^{(k)}, \\
R_{16}^{(k)} &= \delta_0^{(1)} A_{36}^{(k)} + \delta_0^{(3)} A_{66}^{(k)} + c_{44}^{(1)} A_{16}^{(k)} + c_{44}^{(3)} A_{46}^{(k)}, \\
R_{31}^{(k)} &= -a_k \beta_0^{(1)} A_{12}^{(k)} - a_k \beta_0^{(3)} A_{42}^{(k)} + c_{33}^{(1)} A_{13}^{(k)} + c_{33}^{(3)} A_{16}^{(k)}, \\
R_{34}^{(k)} &= -a_k \beta_0^{(1)} A_{42}^{(k)} - a_k \beta_0^{(3)} A_{45}^{(k)} + c_{33}^{(1)} A_{43}^{(k)} + c_{33}^{(3)} A_{46}^{(k)}, \\
R_{43}^{(k)} &= \delta_0^{(4)} A_{33}^{(k)} + \delta_0^{(2)} A_{36}^{(k)} + c_{44}^{(3)} A_{13}^{(k)} + c_{44}^{(2)} A_{43}^{(k)}, \\
R_{46}^{(k)} &= \delta_0^{(4)} A_{36}^{(k)} + \delta_0^{(2)} A_{66}^{(k)} + c_{44}^{(3)} A_{16}^{(k)} + c_{44}^{(2)} A_{46}^{(k)}, \\
R_{61}^{(k)} &= -a_k \beta_0^{(4)} A_{12}^{(k)} - a_k \beta_0^{(2)} A_{42}^{(k)} + c_{33}^{(3)} A_{13}^{(k)} + c_{33}^{(2)} A_{16}^{(k)}, \\
R_{64}^{(k)} &= -a_k \beta_0^{(4)} A_{42}^{(k)} - a_k \beta_0^{(2)} A_{45}^{(k)} + c_{33}^{(3)} A_{43}^{(k)} + c_{33}^{(2)} A_{46}^{(k)}, \quad k = 3, \dots, 6.
\end{aligned} \tag{9}$$

We can easily prove that every column of the matrix $[P(\partial x, n)\Gamma]^*$ is a solution of the system (1) with respect to the point x if $x \neq y$ and all elements $M_{pq}^{(k)}$ have a singularity of type $|x|^{-2}$.

We choose $\delta_0^{(j)}, \beta_0^{(j)}, j = 1, \dots, 4$, so that

$$\begin{aligned}
\sum_{k=3}^6 \frac{R_{13}^{(k)}}{\sqrt{a_k}} = 0, \quad \sum_{k=3}^6 \frac{R_{31}^{(k)}}{\sqrt{a_k}} = 0, \quad \sum_{k=3}^6 \frac{R_{16}^{(k)}}{\sqrt{a_k}} = 0, \quad \sum_{k=3}^6 \frac{R_{34}^{(k)}}{\sqrt{a_k}} = 0, \\
\sum_{k=3}^6 \frac{R_{43}^{(k)}}{\sqrt{a_k}} = 0, \quad \sum_{k=3}^6 \frac{R_{46}^{(k)}}{\sqrt{a_k}} = 0, \quad \sum_{k=3}^6 \frac{R_{61}^{(k)}}{\sqrt{a_k}} = 0, \quad \sum_{k=3}^6 \frac{R_{64}^{(k)}}{\sqrt{a_k}} = 0,
\end{aligned} \tag{10}$$

After some simplification, we find from (10) that

$$\begin{aligned}
\Delta &= \sum_{k=3}^6 A_{12}^{(k)} \sqrt{a_k} \sum_{k=3}^6 A_{45}^{(k)} \sqrt{a_k} - \left(\sum_{k=3}^6 A_{42}^{(k)} \sqrt{a_k} \right)^2 = \\
&= \sqrt{a_3 a_4 a_5 a_6} \left[\sum_{k=3}^6 \frac{A_{33}^{(k)}}{\sqrt{a_k}} \sum_{k=3}^6 \frac{A_{66}^{(k)}}{\sqrt{a_k}} - \left(\sum_{k=3}^6 \frac{A_{36}^{(k)}}{\sqrt{a_k}} \right)^2 \right] = \\
&= \frac{B_0}{b_0^2} [(\delta_{11} \delta_{22} + b_0 m_1 m_3) q_4 + q_1 b_4 + \delta_{22} b_0 m_2] (\sqrt{a_3 a_4 a_5 a_6})^{-1} + \\
&\quad + q_1 (\delta_{11} \delta_{22} + b_0 m_1 m_3 - k_1) + b_0 \delta_{11} m_2],
\end{aligned}$$

where

$$\begin{aligned}
q_1 &= c_{11}^{(1)} c_{11}^{(2)} - c_{11}^{(3)2}, \quad q_4 = c_{44}^{(1)} c_{44}^{(2)} - c_{44}^{(3)2}, \quad b_0 = q_1 q_4, \\
m_1 &= \sum_{k=3}^6 \sqrt{a_k}, \quad m_2 = \sum_{p \neq q} \sqrt{a_p a_q}, \\
m_3 &= \sum_{p \neq q \neq j} \sqrt{a_p a_q a_j}, \quad p, q, j = 3, \dots, 6,
\end{aligned}$$

$$\begin{aligned}
\delta_{11} &= c_{11}^{(1)} c_{44}^{(2)} + c_{44}^{(1)} c_{11}^{(2)} - 2c_{11}^{(3)} c_{44}^{(3)} > 0, \\
\delta_{22} &= c_{33}^{(1)} c_{44}^{(2)} + c_{44}^{(1)} c_{33}^{(2)} - 2c_{33}^{(3)} c_{44}^{(3)} > 0, \\
k_1 + k_2 &= 2(\alpha_{13}^{(1)} \alpha_{13}^{(2)} - \alpha_{13}^{(3)2}) - \alpha_{13}^{(1)} v_{11} - \alpha_{13}^{(2)} w_{14} - \alpha_{13}^{(3)}(w_{12} + v_{21}), \\
k_1 &= \frac{1}{c_{44}^{(2)2}} \left[c_{44}^{(2)2} c_{13}^{(3)} - 2c_{44}^{(2)} c_{44}^{(3)} c_{13}^{(3)} + c_{44}^{(3)2} c_{13}^{(2)} + c_{44}^{(2)} \right]^2 + \\
&\quad + \frac{2q_4}{c_{44}^{(2)2}} \left[c_{44}^{(2)} c_{13}^{(3)} - c_{44}^{(3)} c_{13}^{(2)} \right]^2 + \frac{q_4^2}{c_{44}^{(2)2}} \alpha_{13}^{(2)2}, \\
B_0^{-1} &= \prod_{p \neq q} (\sqrt{a_p} + \sqrt{a_q}), \quad p, q = 3, \dots, 6.
\end{aligned}$$

Taking into account the inequalities obtained from the positive definiteness of the energy $E(u, u)$, we conclude that $\Delta \neq 0$. When $\delta_0^{(j)}, \beta_0^{(j)}$ are solutions of the system (10), we denote the vector $P(\partial y, n)U$, by $N(\partial y, n)U$. Then from (4), when $U^+ = 0$, we have

$$U(x) = -\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Gamma(y-x) N(\partial y, n) U dy_1 dy_2.$$

Hence for the vector NU as $x(x_1, x_2, x_3) \rightarrow z(z_1, z_2, 0)$ we find

$$[N(\partial z, n)U]^+ + \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} N\Gamma(y-z)(NU)^+ dy_1 dy_2 = 0.$$

Note that $N\Gamma(z-y) = 0$, $z \in S$. Therefore $(NU)^+ = 0$, and from (4) we have $U = 0, x \in D$. Therefore the homogeneous equation has only the trivial solution. Thus we formulate the following

Theorem. *The first BVP has at most one regular solution.*

The first BVP. A solution of the first BVP will be sought in the domain D in terms of the double layer potential

$$U(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} [N(\partial y, n)\Gamma(y-x)]^* g(y) dy_1 dy_2, \quad (11)$$

where g is an unknown real vector. Taking into account the properties of the double layer potential and the boundary condition for determining g , we obtain the following Fredholm integral equation of second kind:

$$g(z) + \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} [N(\partial y, n)\Gamma(y-z)]^* g(y) dy_1 dy_2 = f(z),$$

Taking into account the fact that $[N\Gamma]^* = 0$, $x_3 = 0$, from the latter equation we have $g(z) = f(z)$ and (11) takes the form

$$U(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} [N(\partial y, n)\Gamma(y-x)]^* f(y) dy_1 dy_2. \quad (12)$$

Thus we have obtained the Poisson formula for the solution of the first BVP for the half-space. Note that (12) is valid if and only if $f \in C^{1,\alpha}(S)$ and satisfies the condition $f = O(\frac{A}{|x|^{1+\beta}})$ at infinity, where A is a constant vector and $\beta > 0$.

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Author's address:

I. Javakhishvili Tbilisi State University
2, University St., Tbilisi 0186
Georgia
E-mail: lbits@viam.sci.tsu.ge